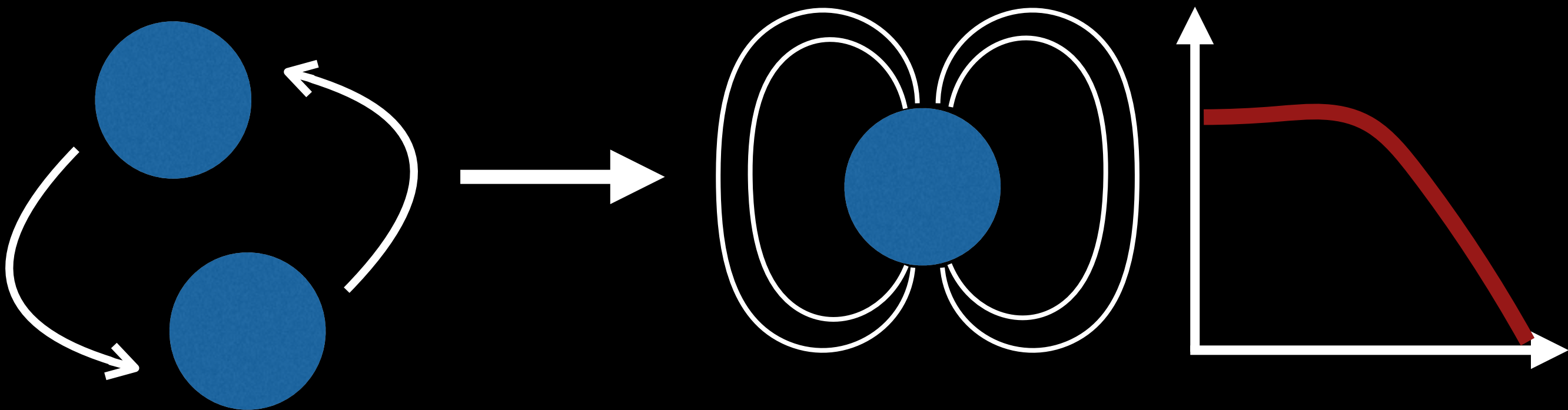
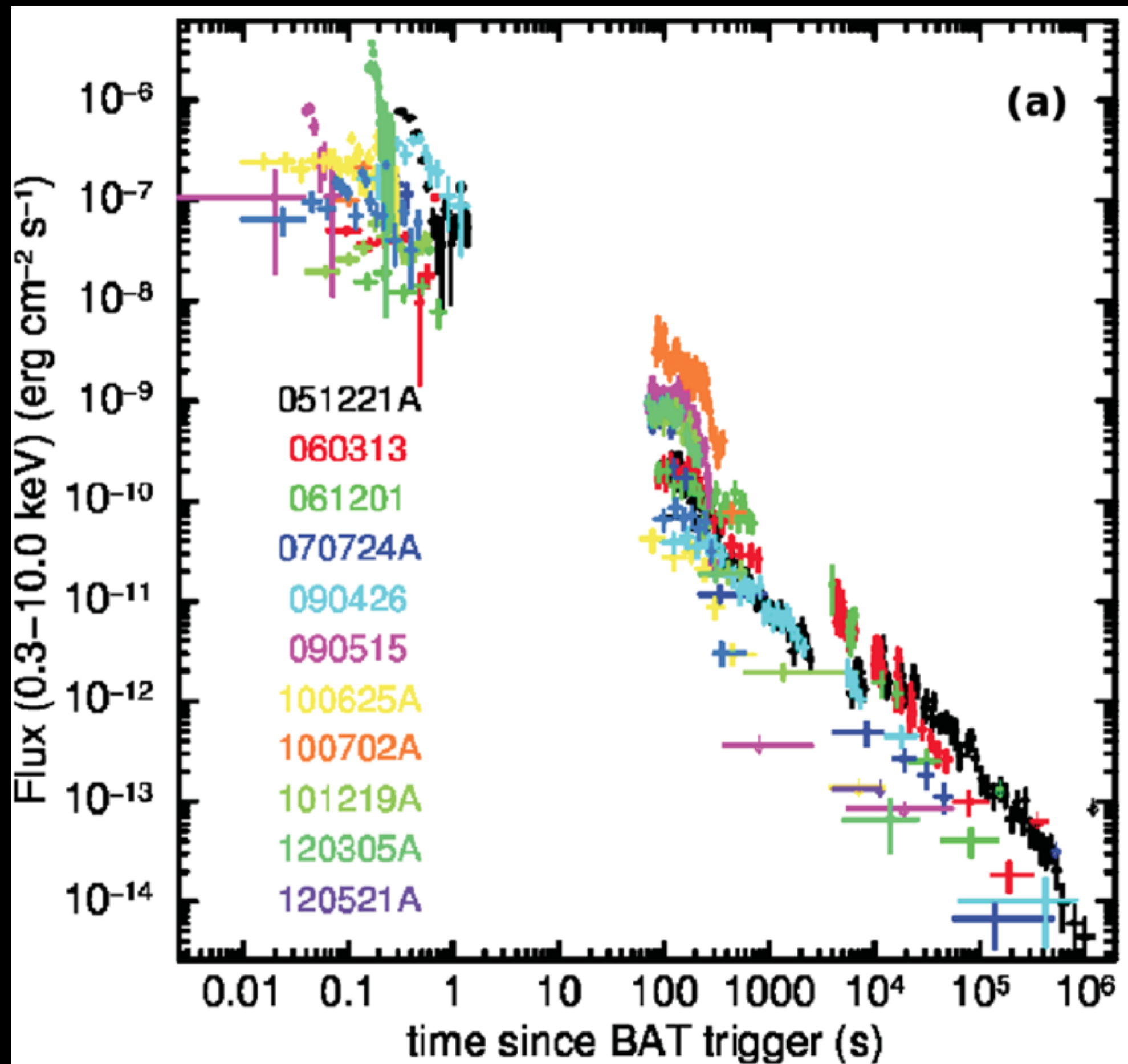
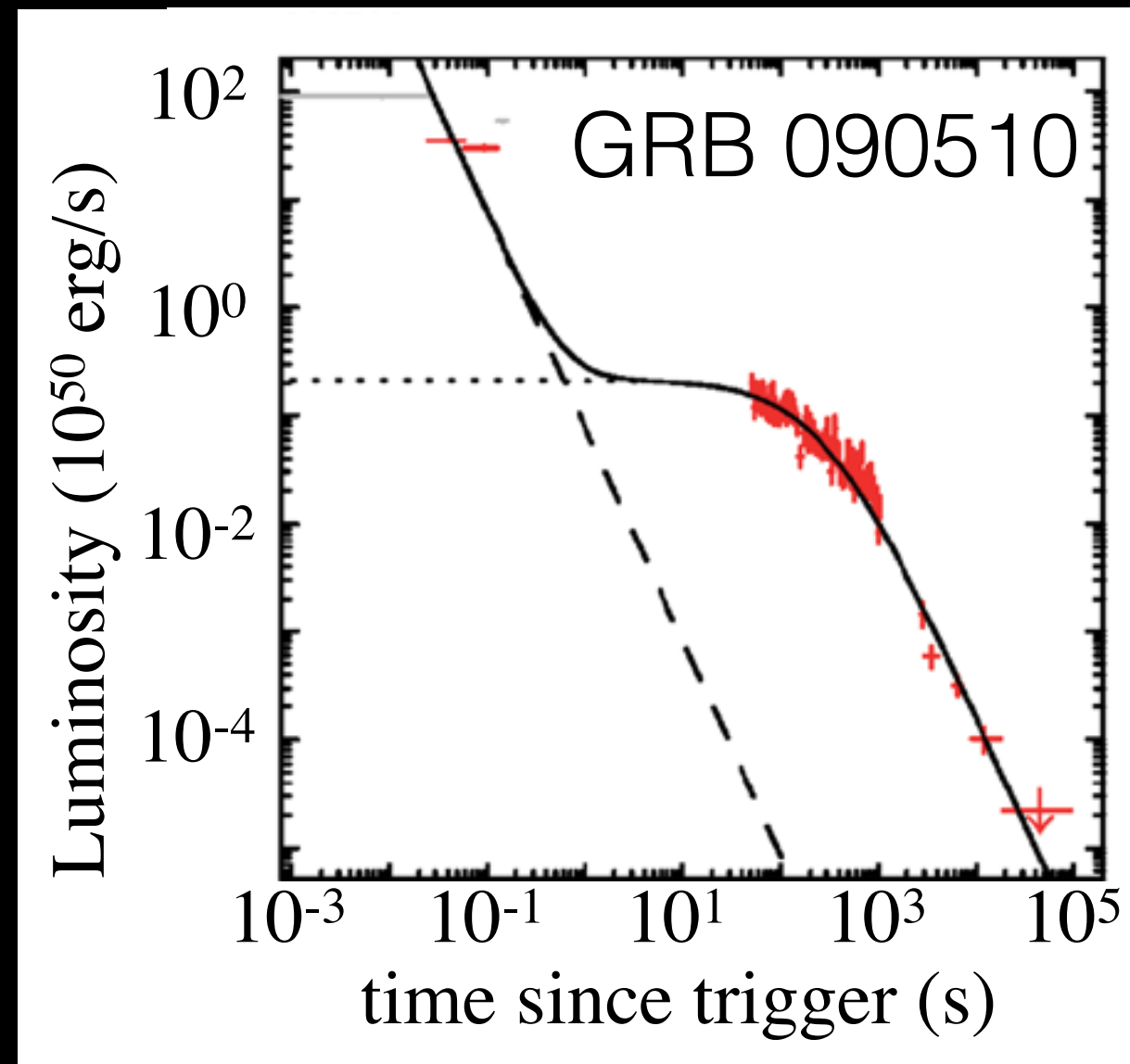
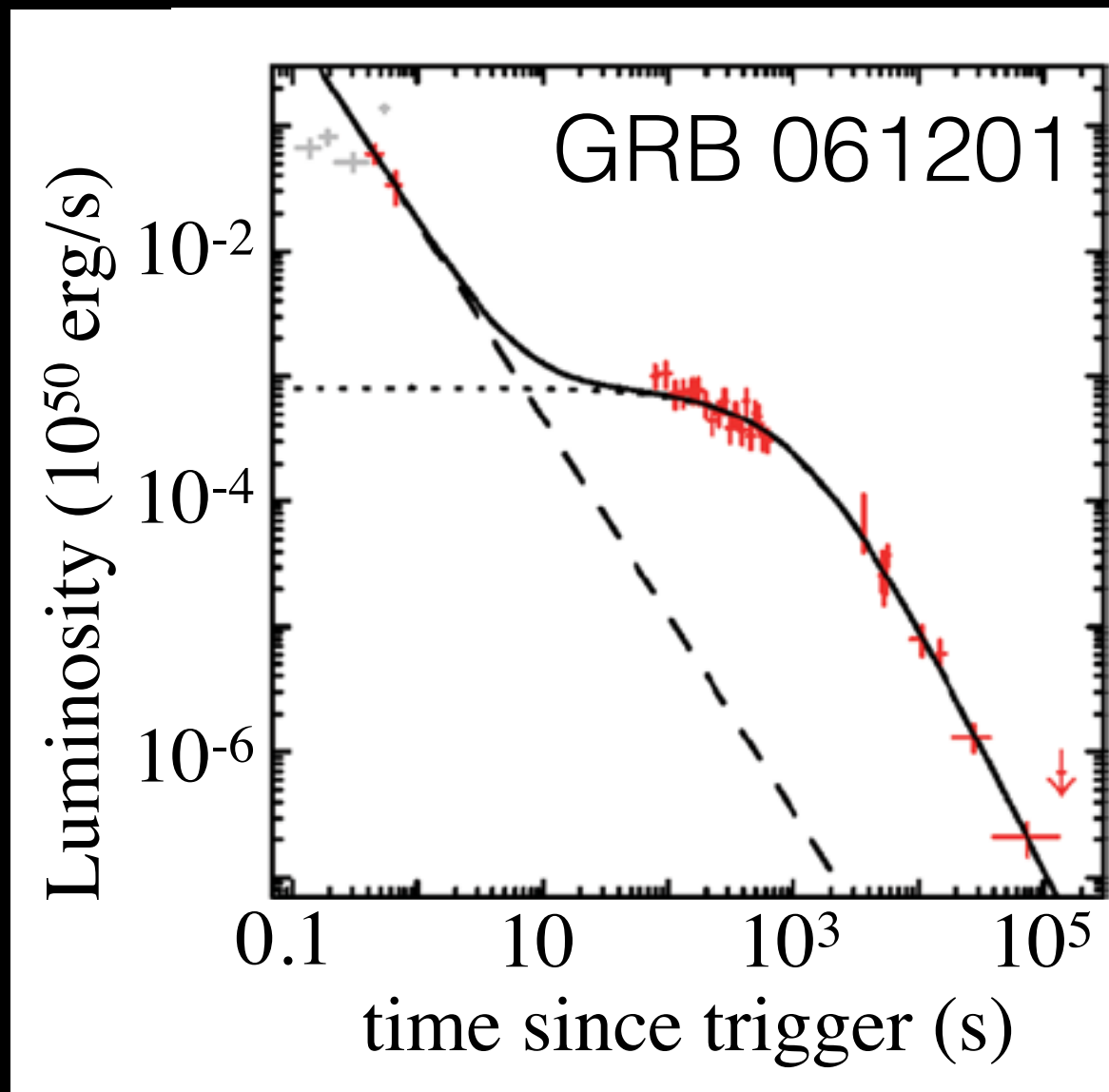


Equation of state and gravitational waves from short GRB remnants

Paul Lasky





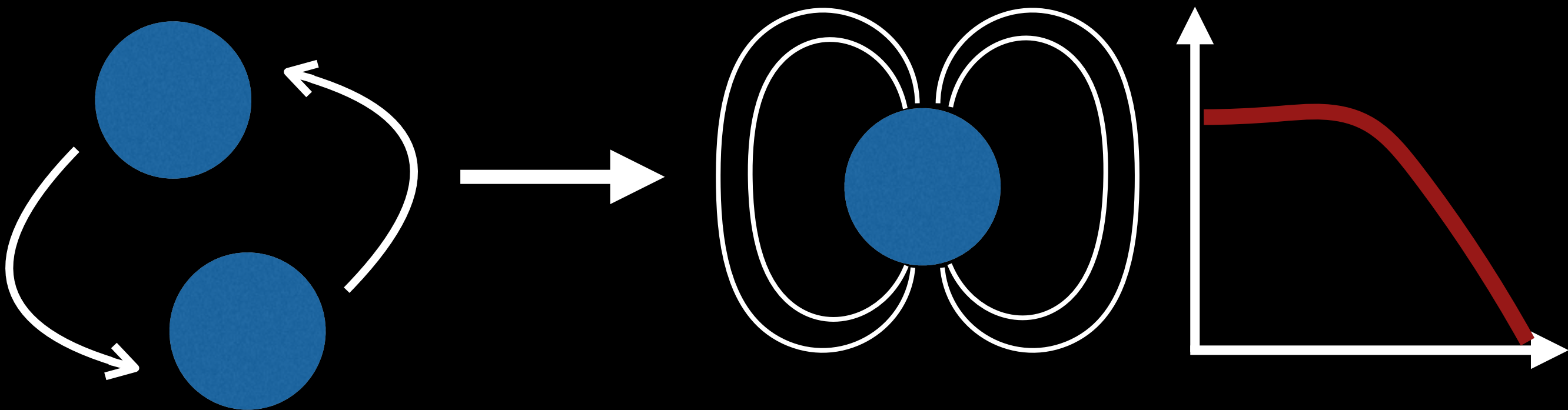


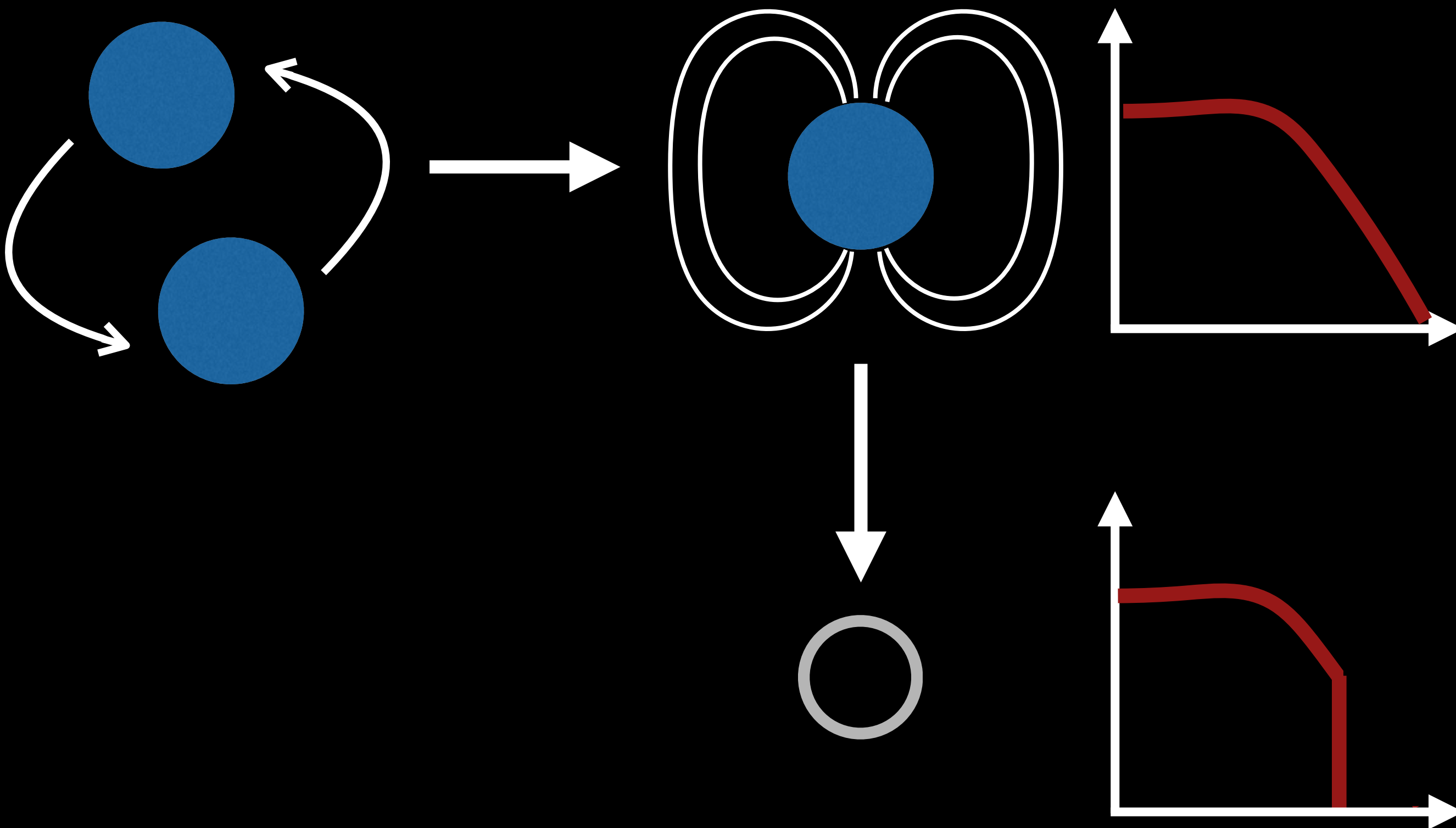
$$B_p^2 \sim \frac{I^2}{R^6 L_0 T_{\text{em}}}$$

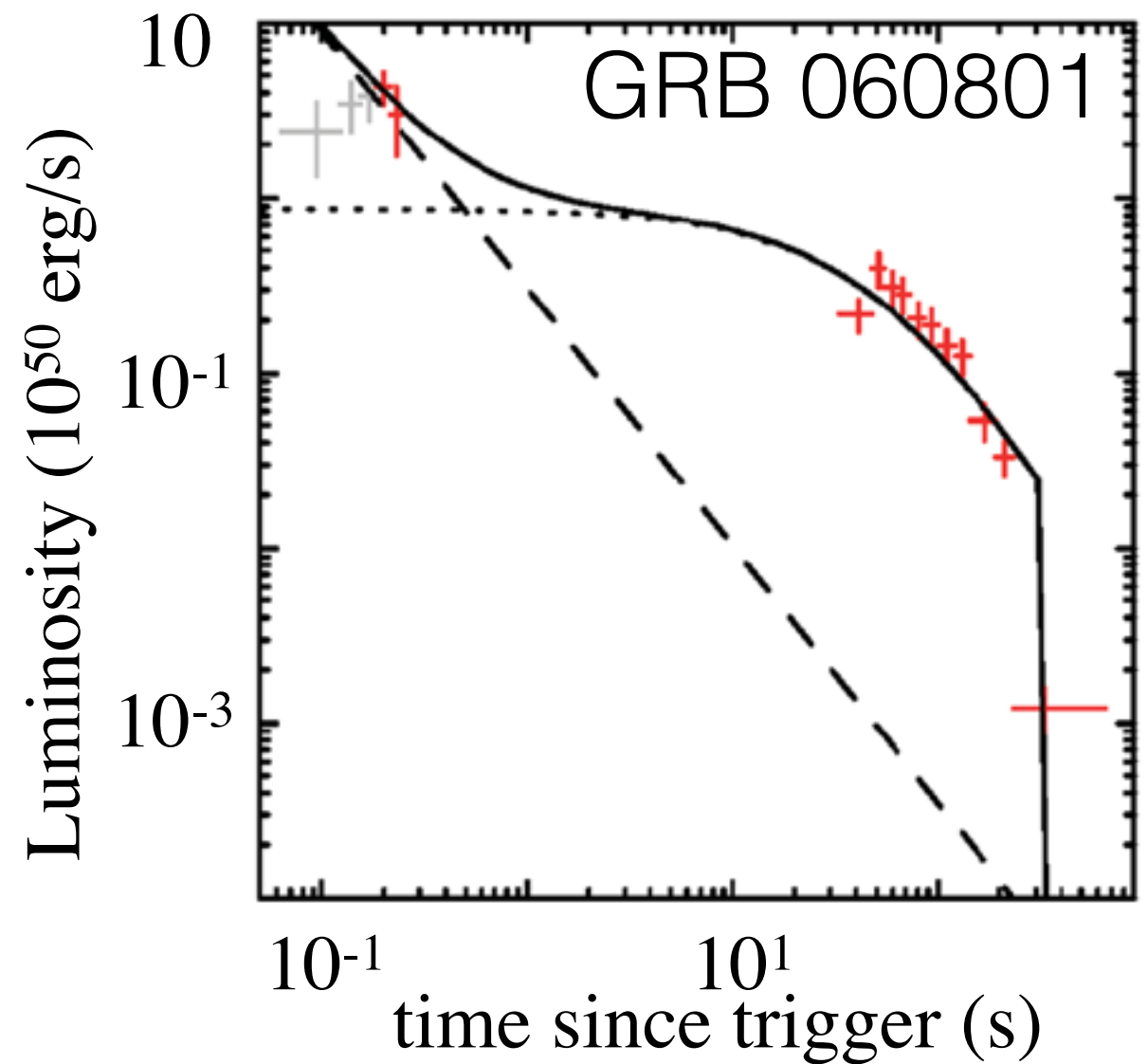
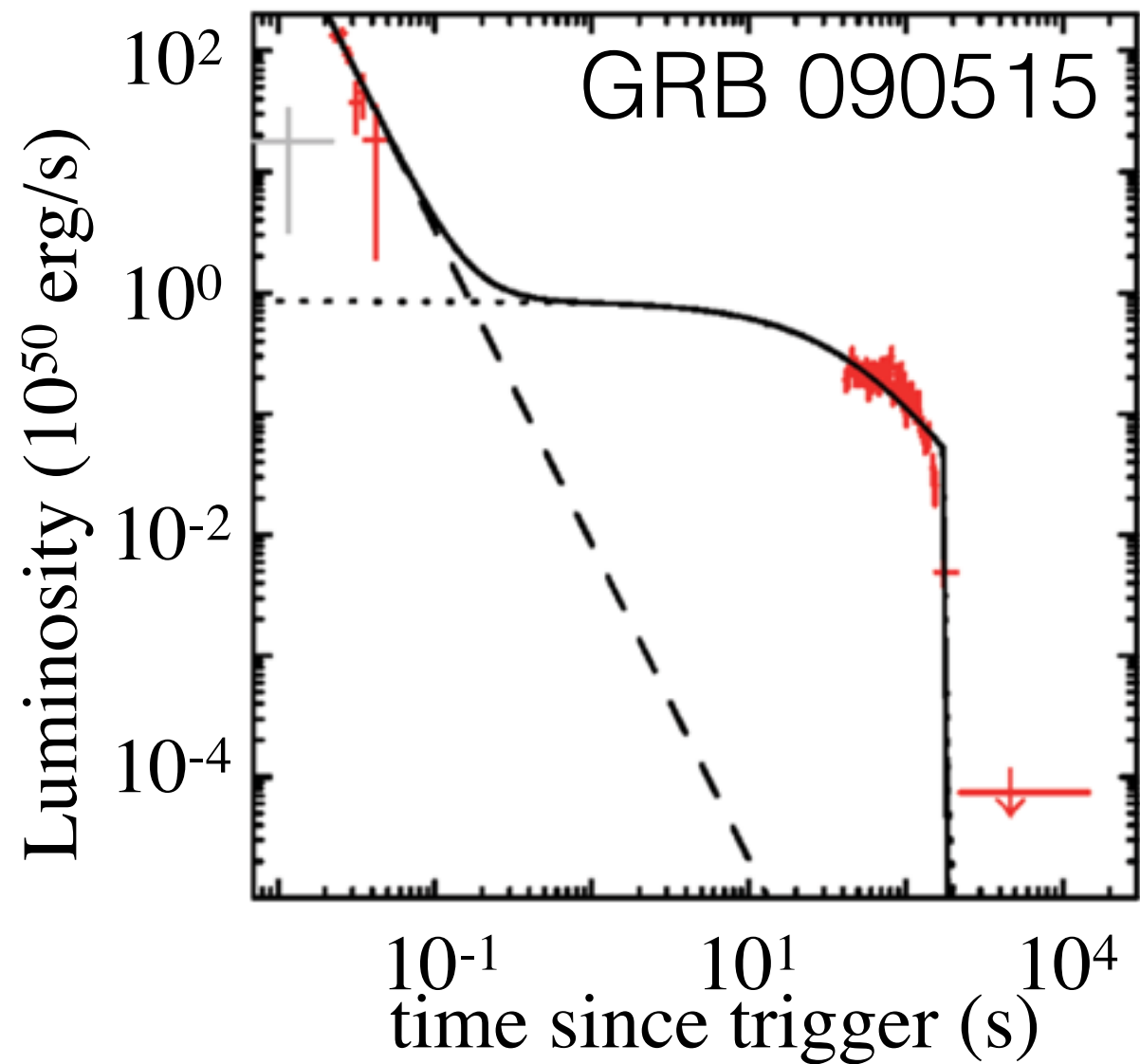
$$P_0^2 \sim \frac{I}{L_0 T_{\text{em}}}$$

L_0 - plateau luminosity

T_{em} - plateau duration



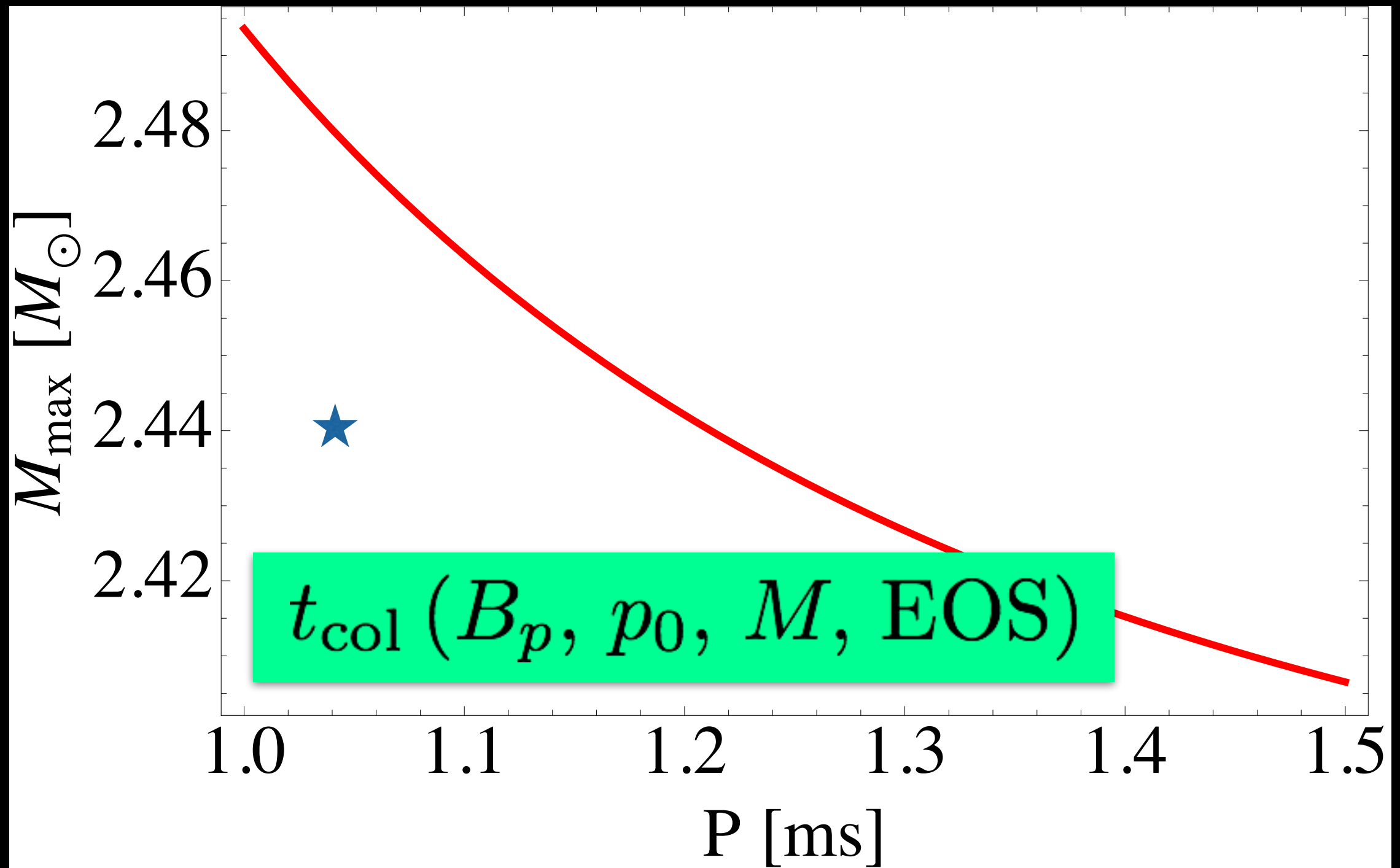




GRB	z	p_0 [ms]	B_p [10^{15} G]	t_{col} [s]
060801	1.13	$1.95^{+0.15}_{-0.13}$	$11.24^{+1.93}_{-1.78}$	326
070724A	0.46	$1.80^{+1.04}_{-0.38}$	$28.72^{+1.42}_{-1.29}$	90
080905A	0.122	$9.80^{+0.78}_{-0.77}$	$39.26^{+10.24}_{-12.16}$	274
101219A	0.718	$0.95^{+0.05}_{-0.05}$	$2.81^{+0.47}_{-0.39}$	138
051221A	0.55	$7.79^{+0.31}_{-0.28}$	$1.80^{+0.14}_{-0.13}$	—
070809	0.219	$5.54^{+0.48}_{-0.43}$	$2.06^{+0.48}_{-0.42}$	—
090426	2.6	$1.89^{+0.08}_{-0.07}$	$4.88^{+0.88}_{-0.90}$	—
090510	0.9	$1.86^{+0.04}_{-0.03}$	$5.06^{+0.27}_{-0.23}$	—

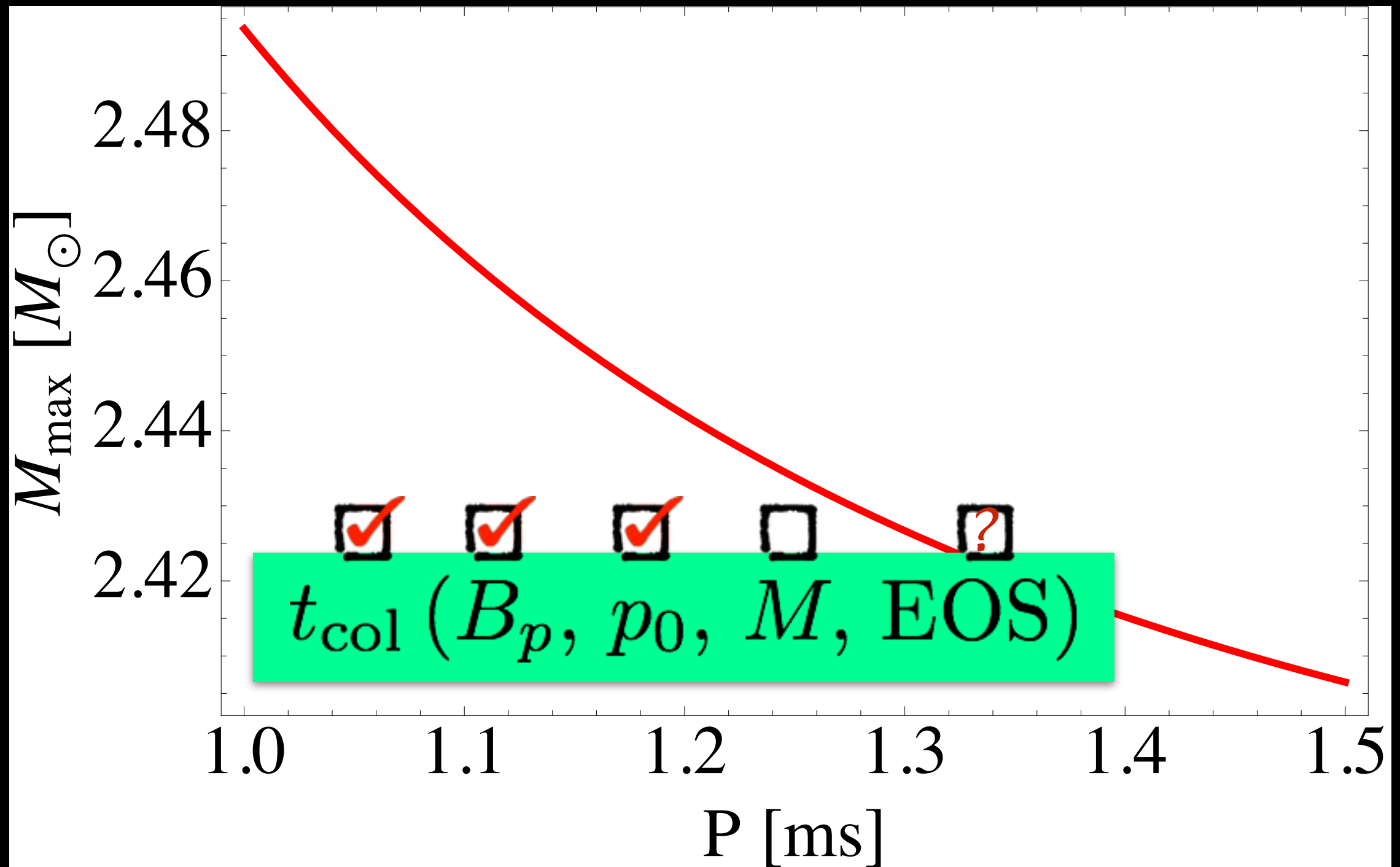
Supramassive Stars

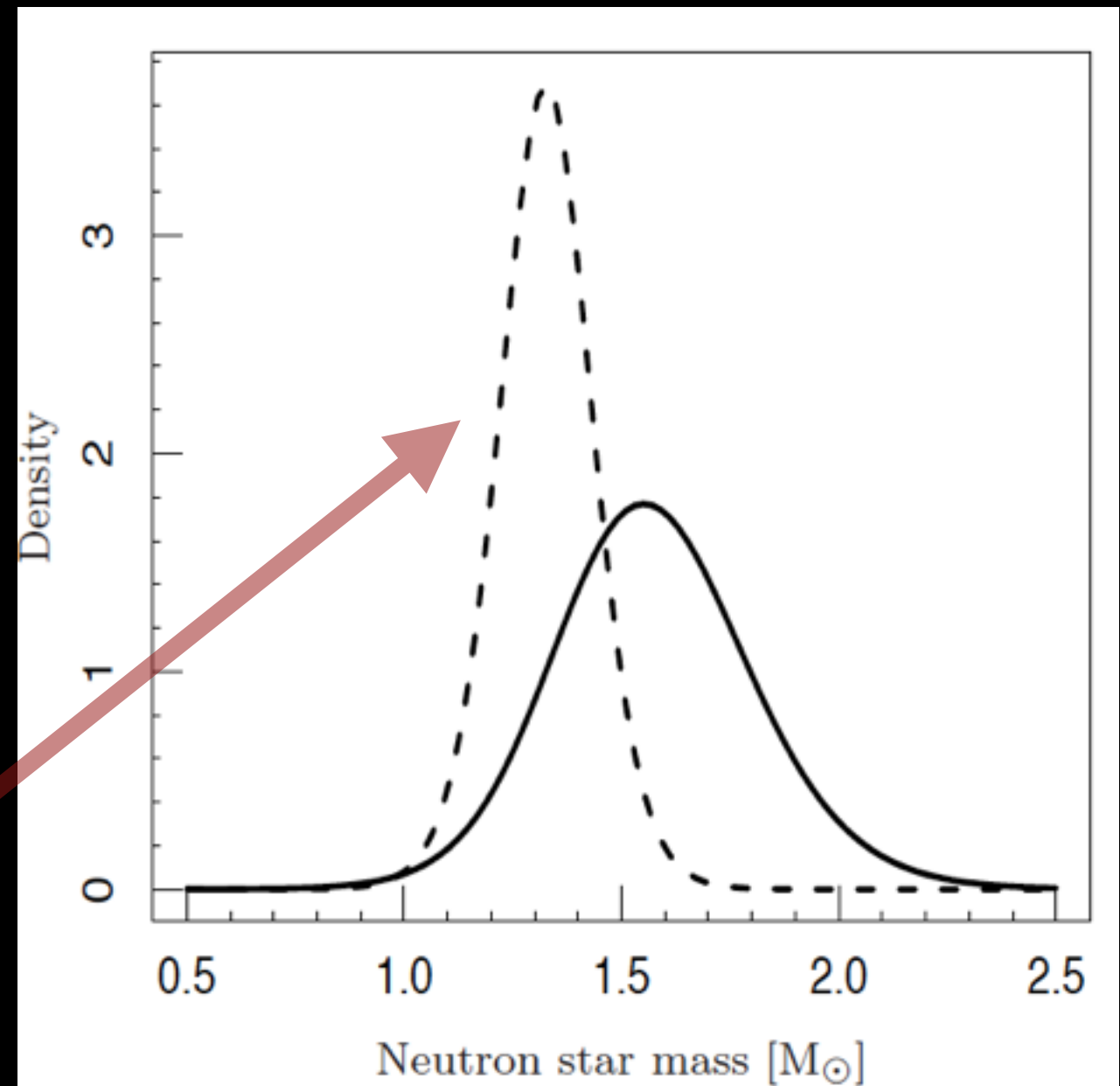
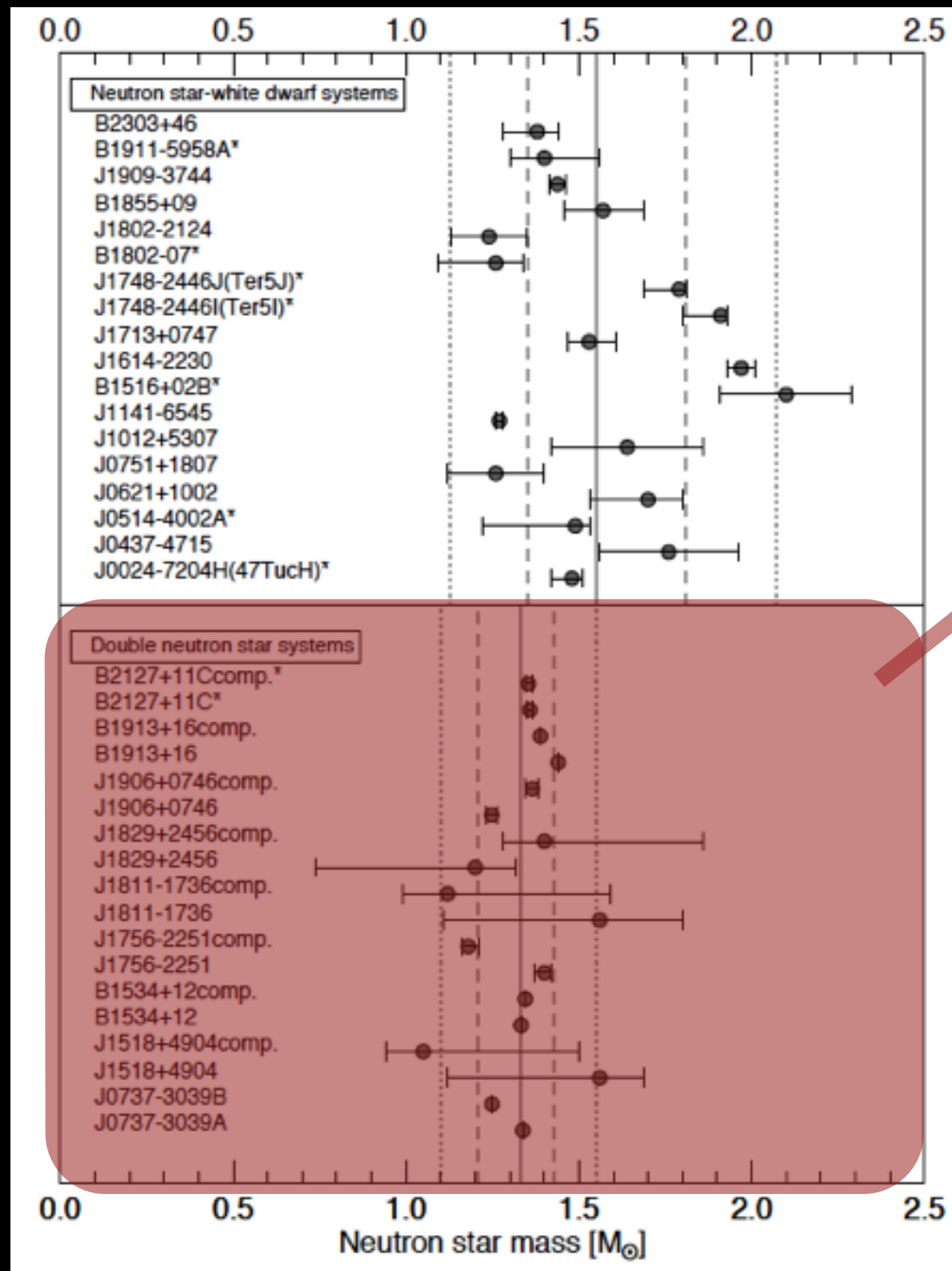
Consider a star: $p_0 = 1.05$ ms
 $M = 2.44 M_{\text{sun}}$



Supramassive Stars

Consider a star: $p_0 = 1.05$ ms
 $M = 2.44 M_{\text{sun}}$

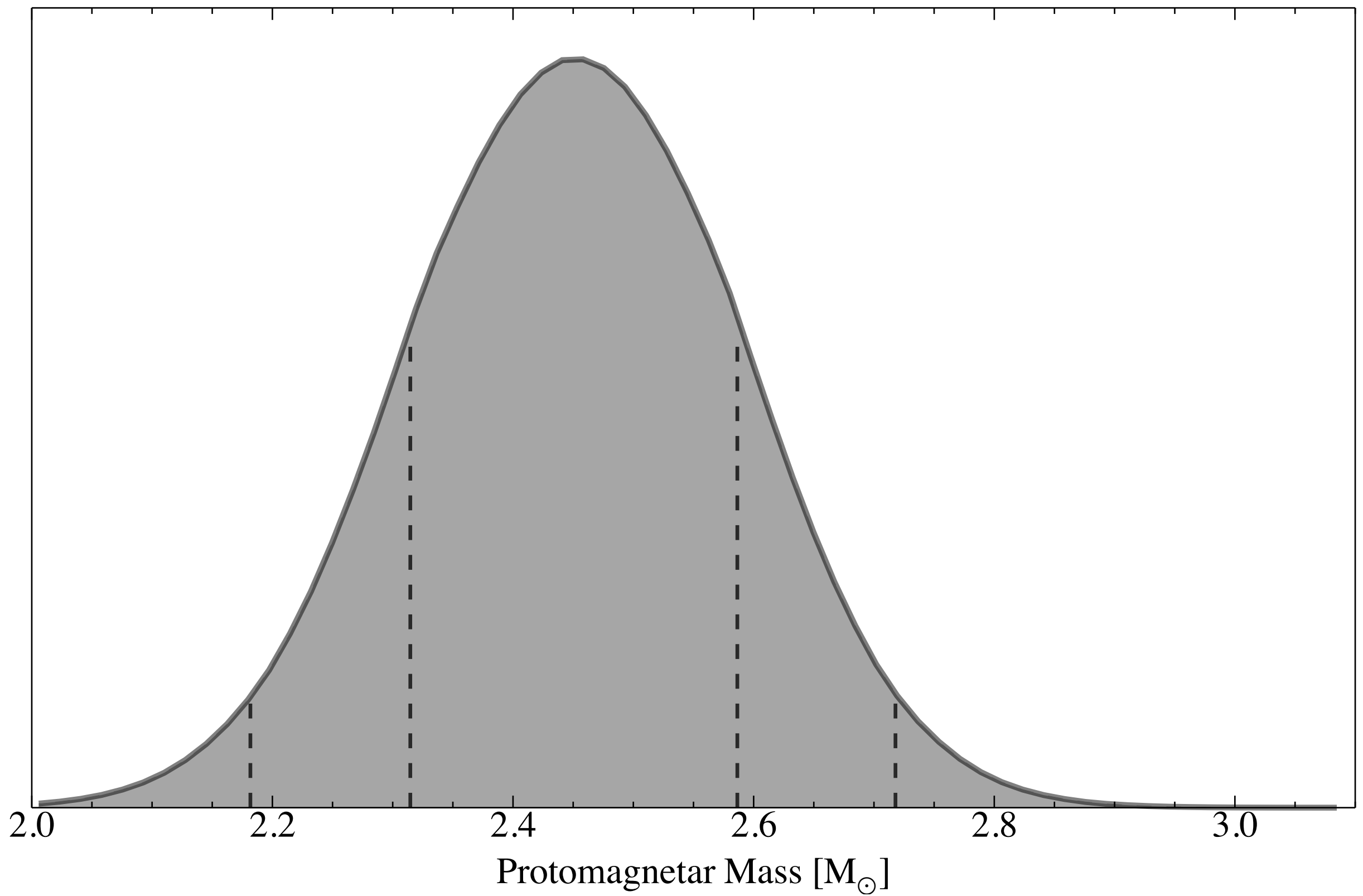




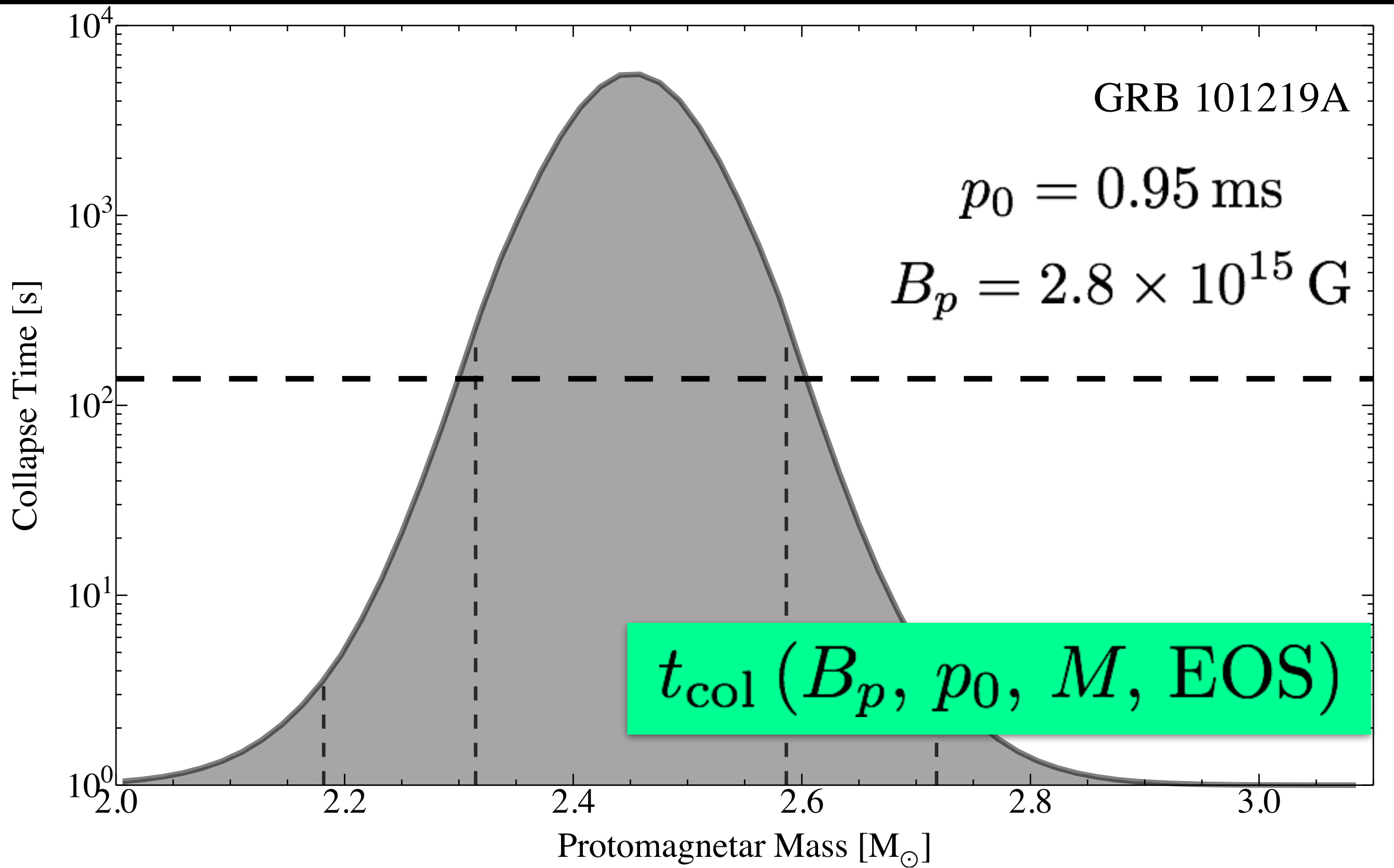
$$M = 1.32^{+0.11}_{-0.11} M_{\odot}$$

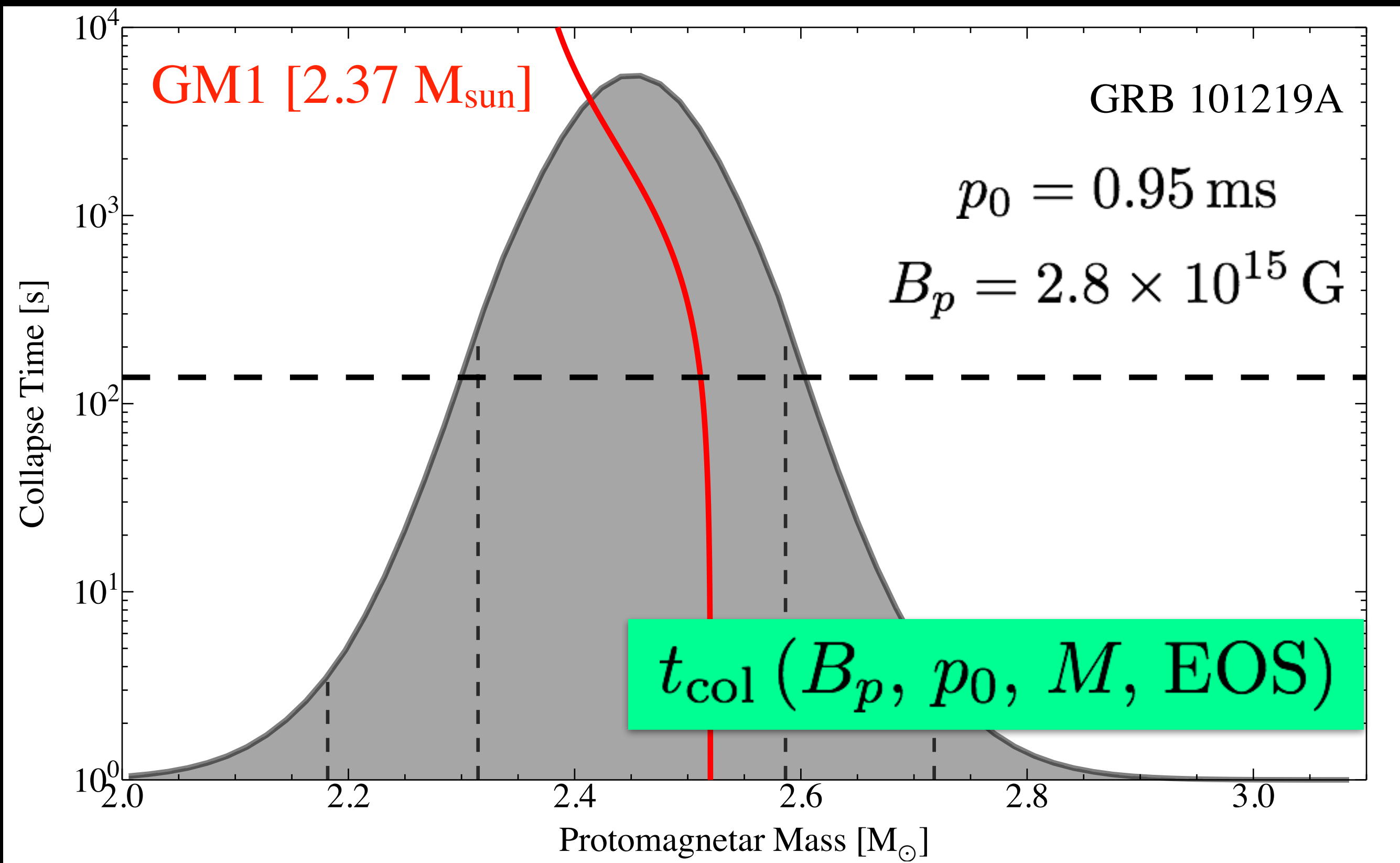
Rest mass conservation

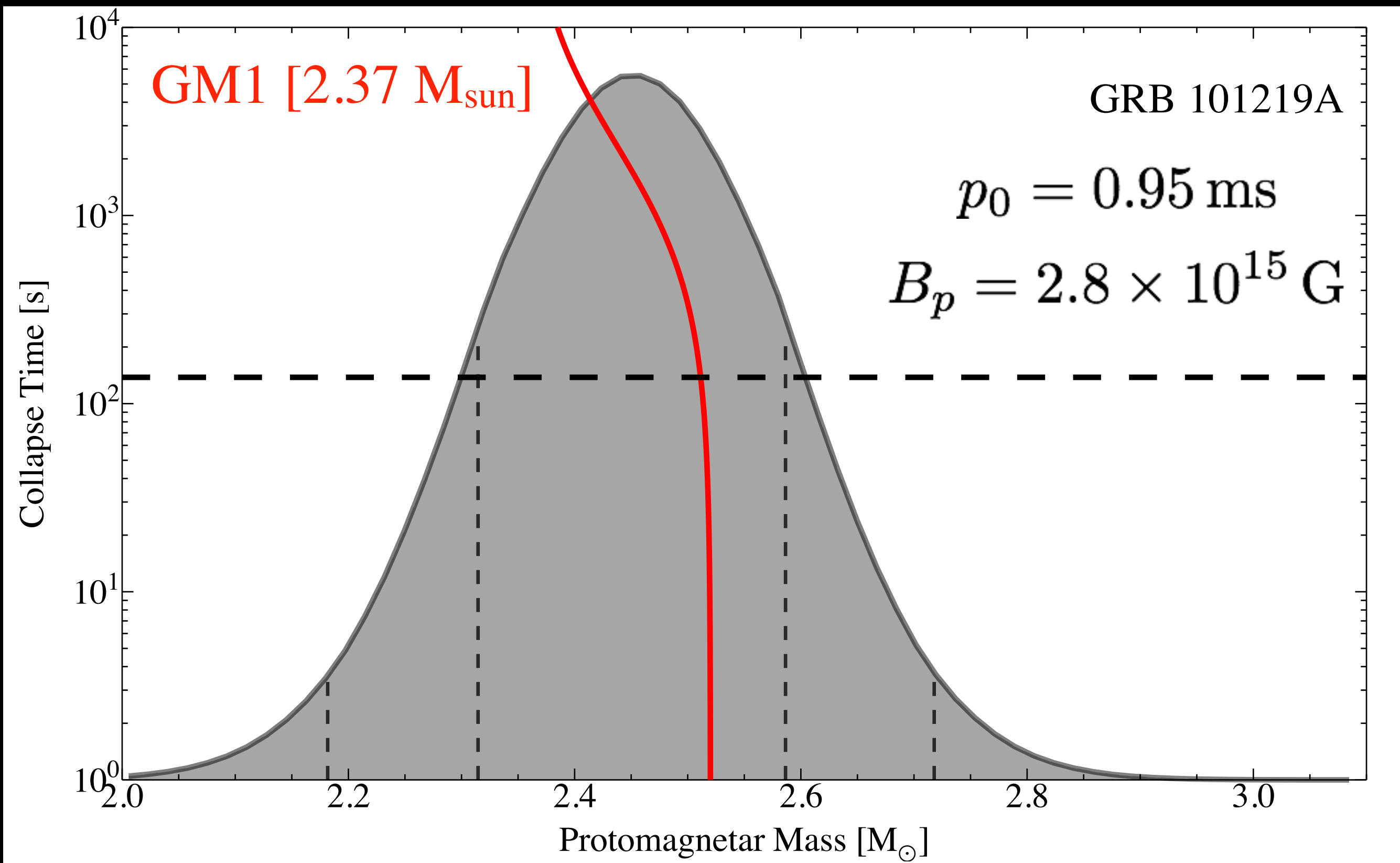
$$M_p = 2.46^{+0.13}_{-0.15} M_{\odot}$$

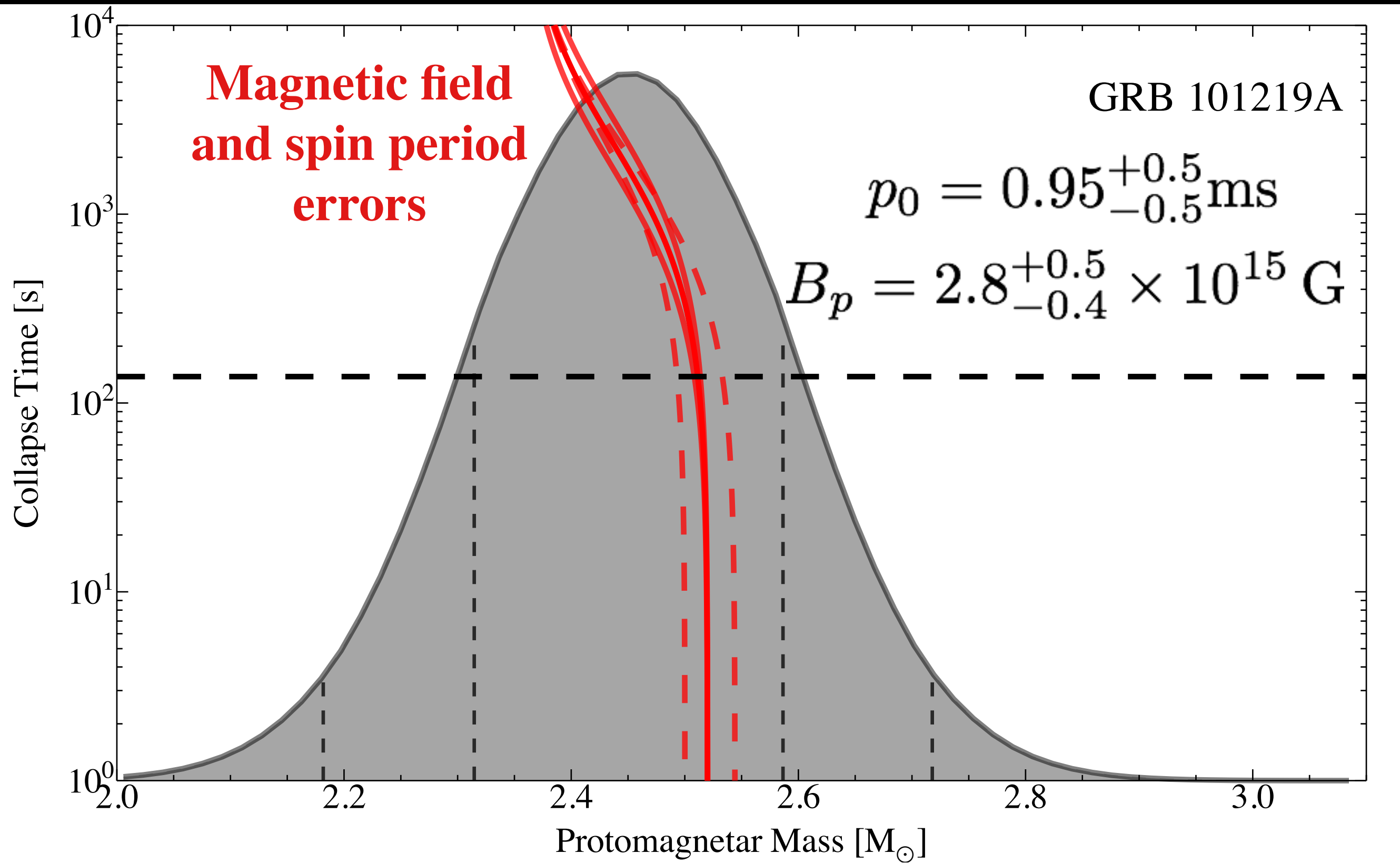


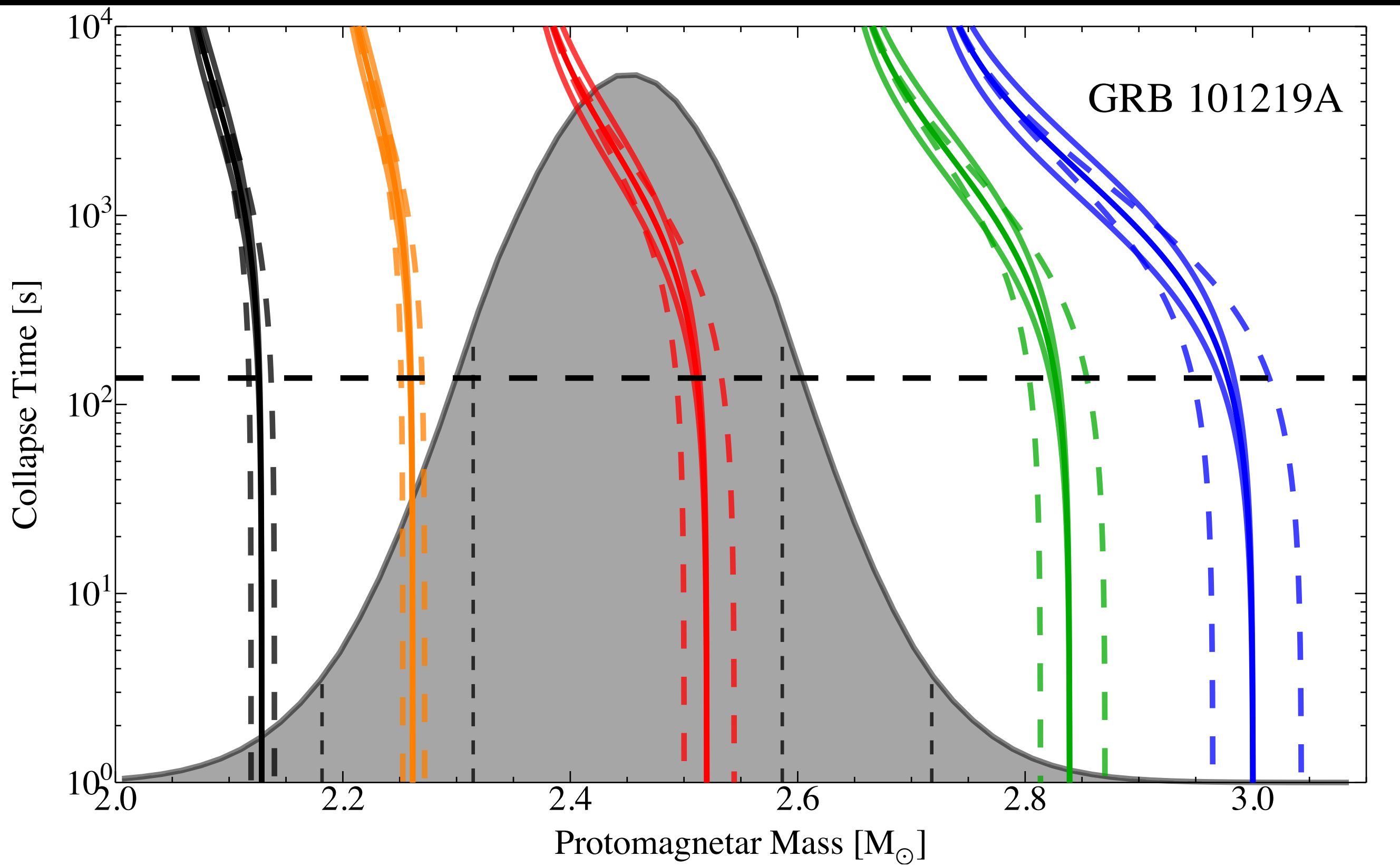
PL, Haskell, Ravi, Howell & Coward (2014)

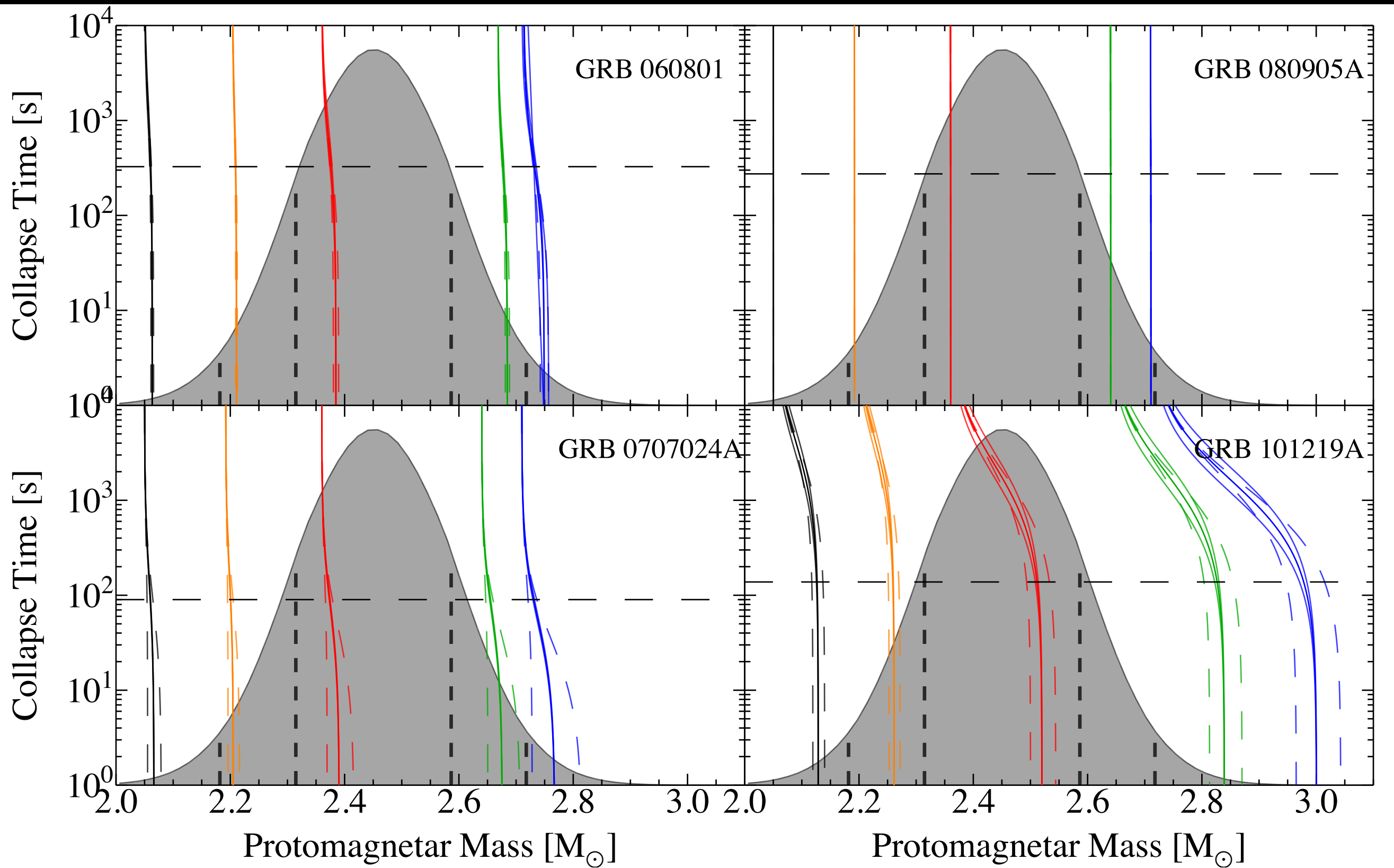






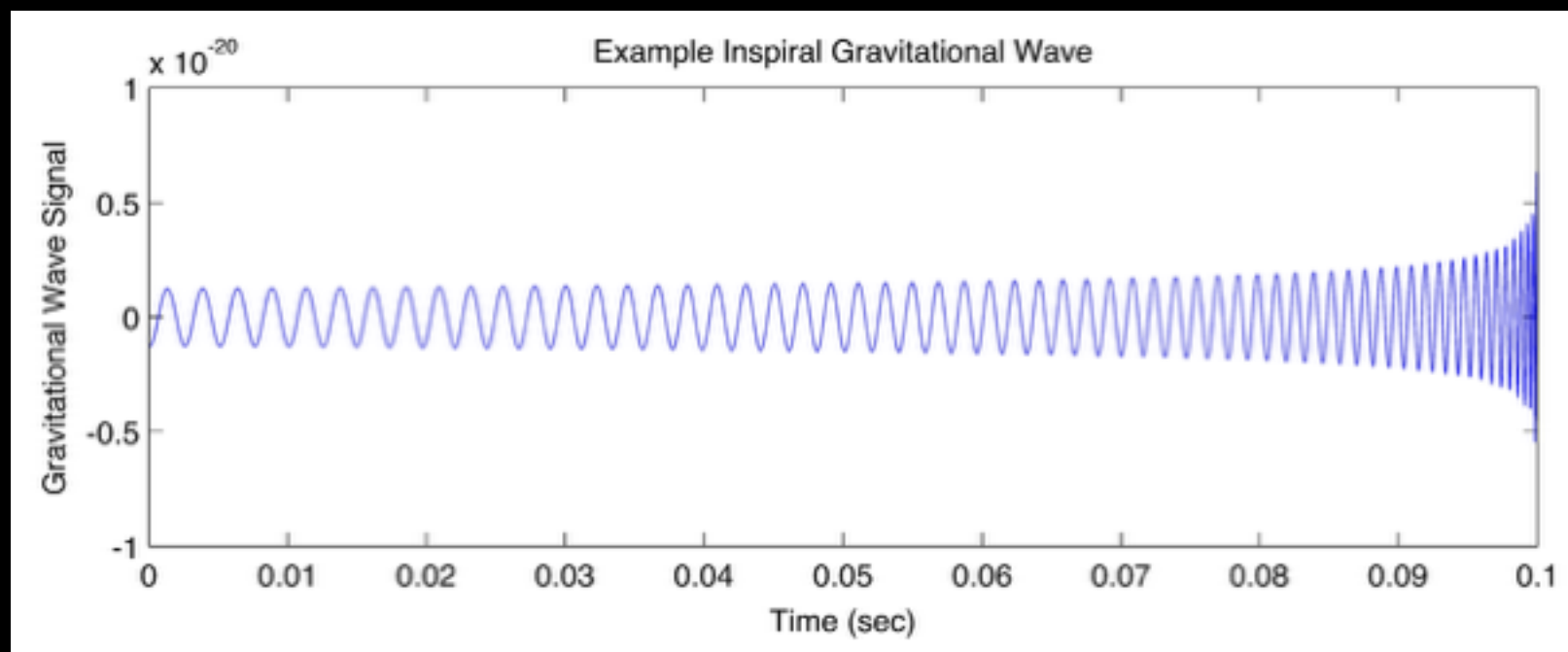
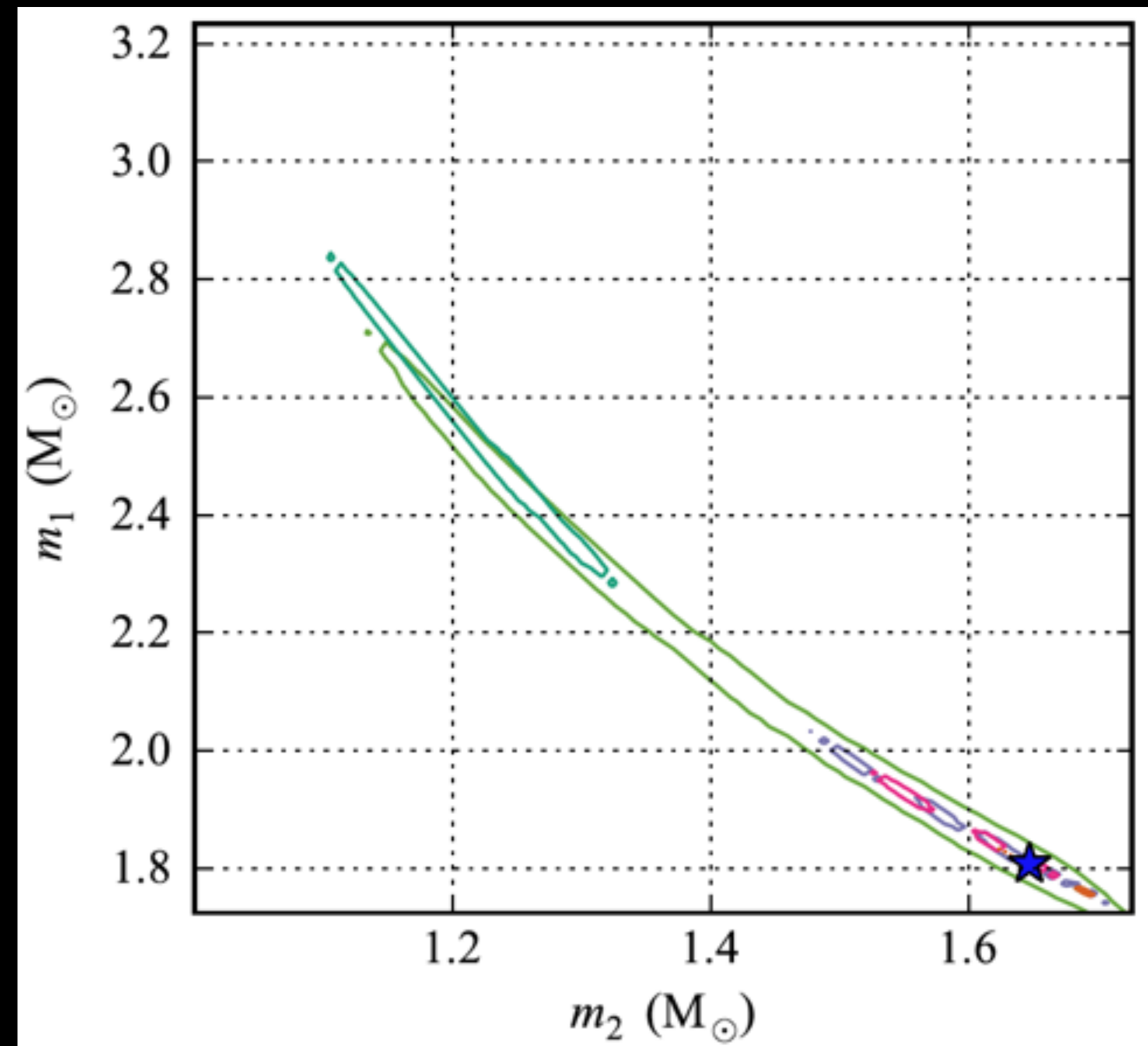






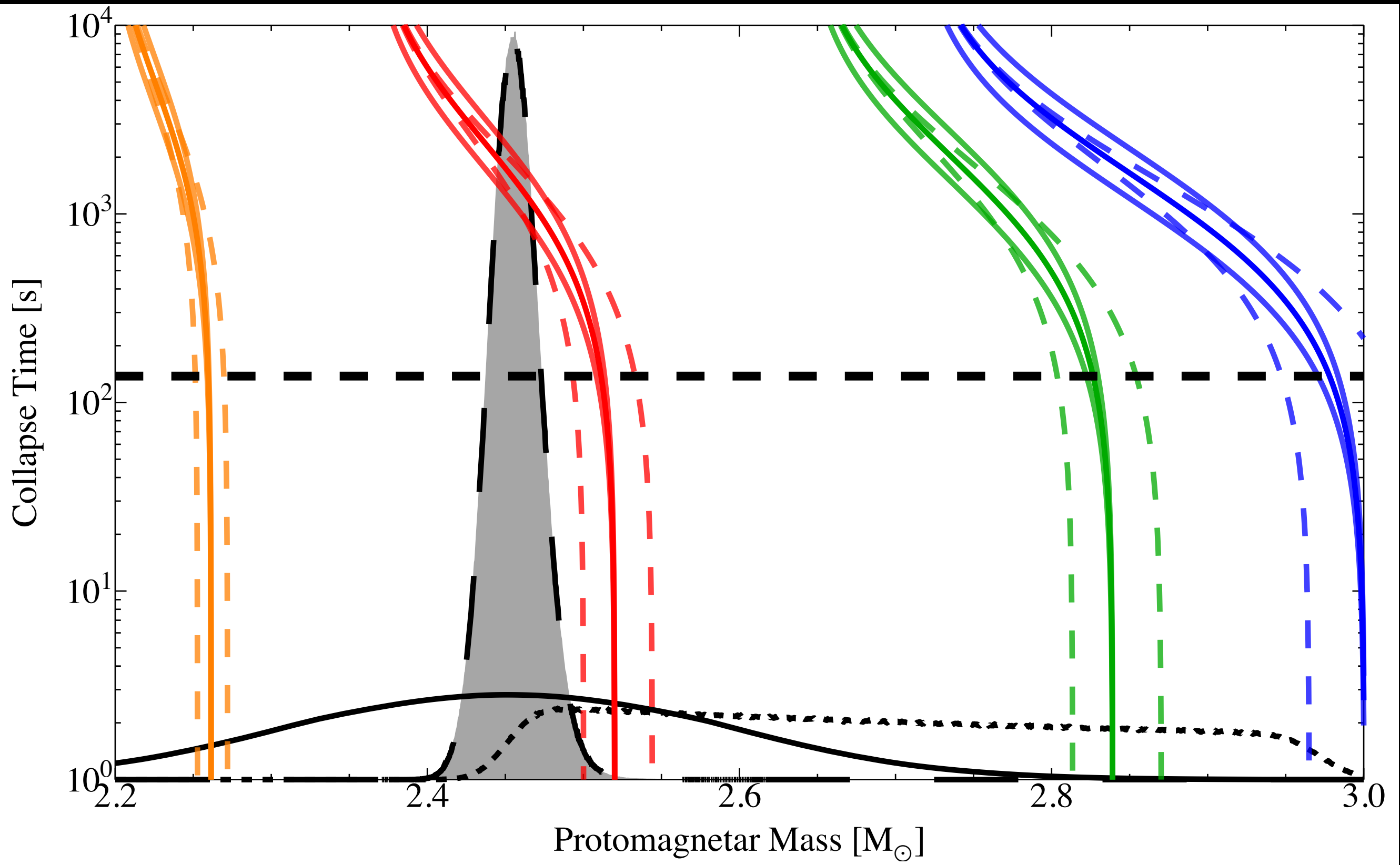
Gravitational Waves

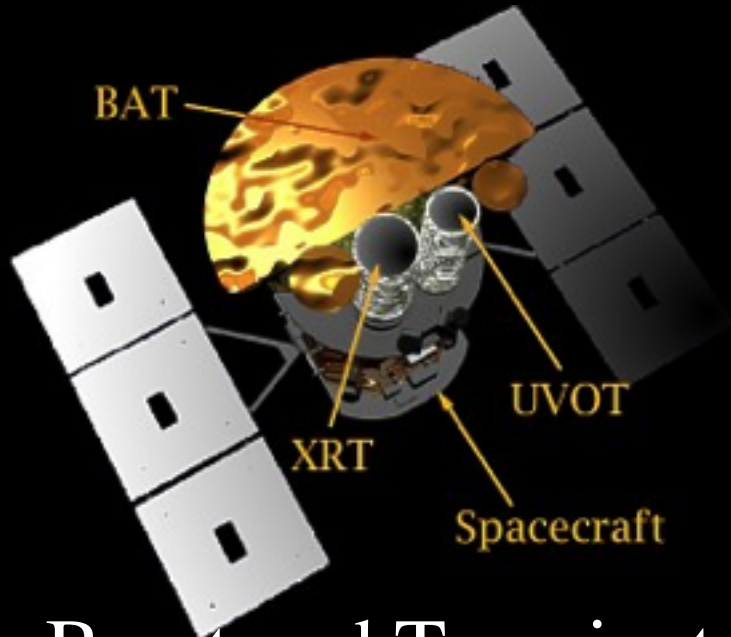
Aasi et al. (2013)



$$\mathcal{M} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{-1/5}}$$

$$\eta = \frac{m_1 m_2}{(m_1 + m_2)^2}$$



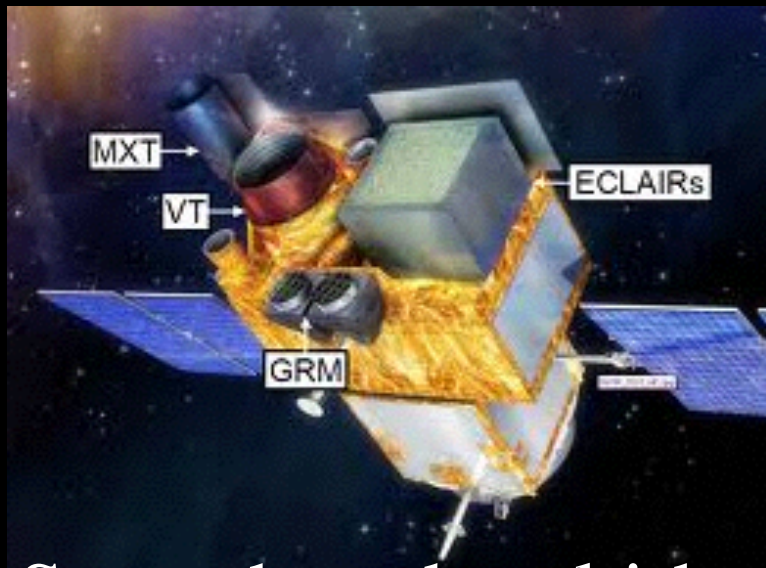


+



$\approx 0.2 \text{ yr}^{-1}$

Burst and Transient Source Experiment (BATSE-XRT)

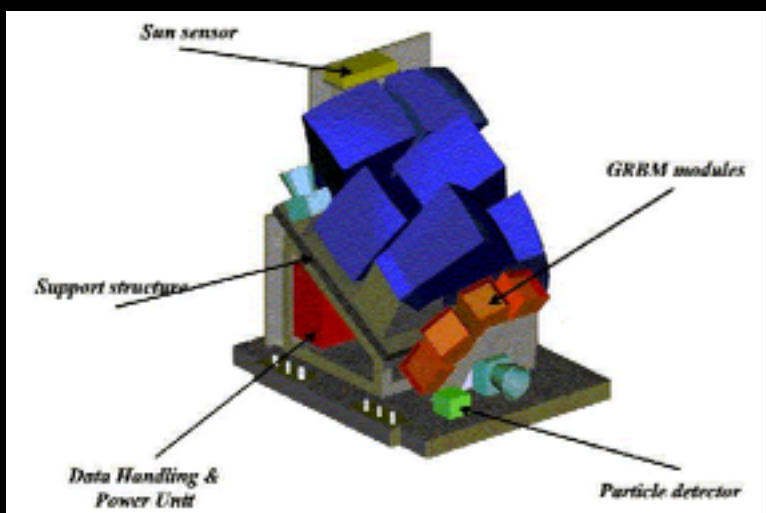


+



$\approx 0.4 \text{ yr}^{-1}$

Space-based multi-band astronomical Variable Object Monitor (SVOM)



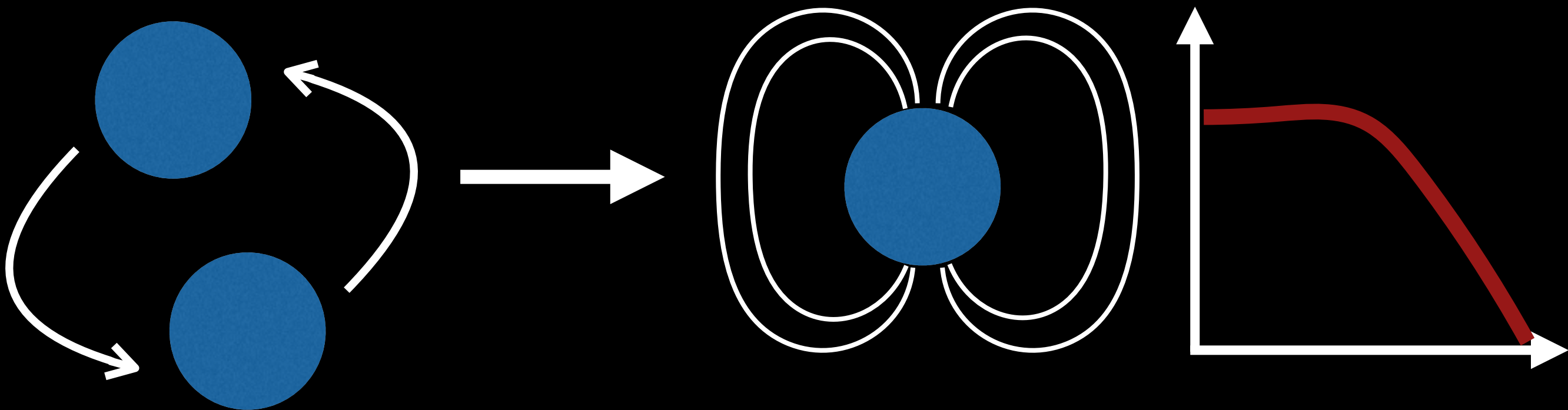
+

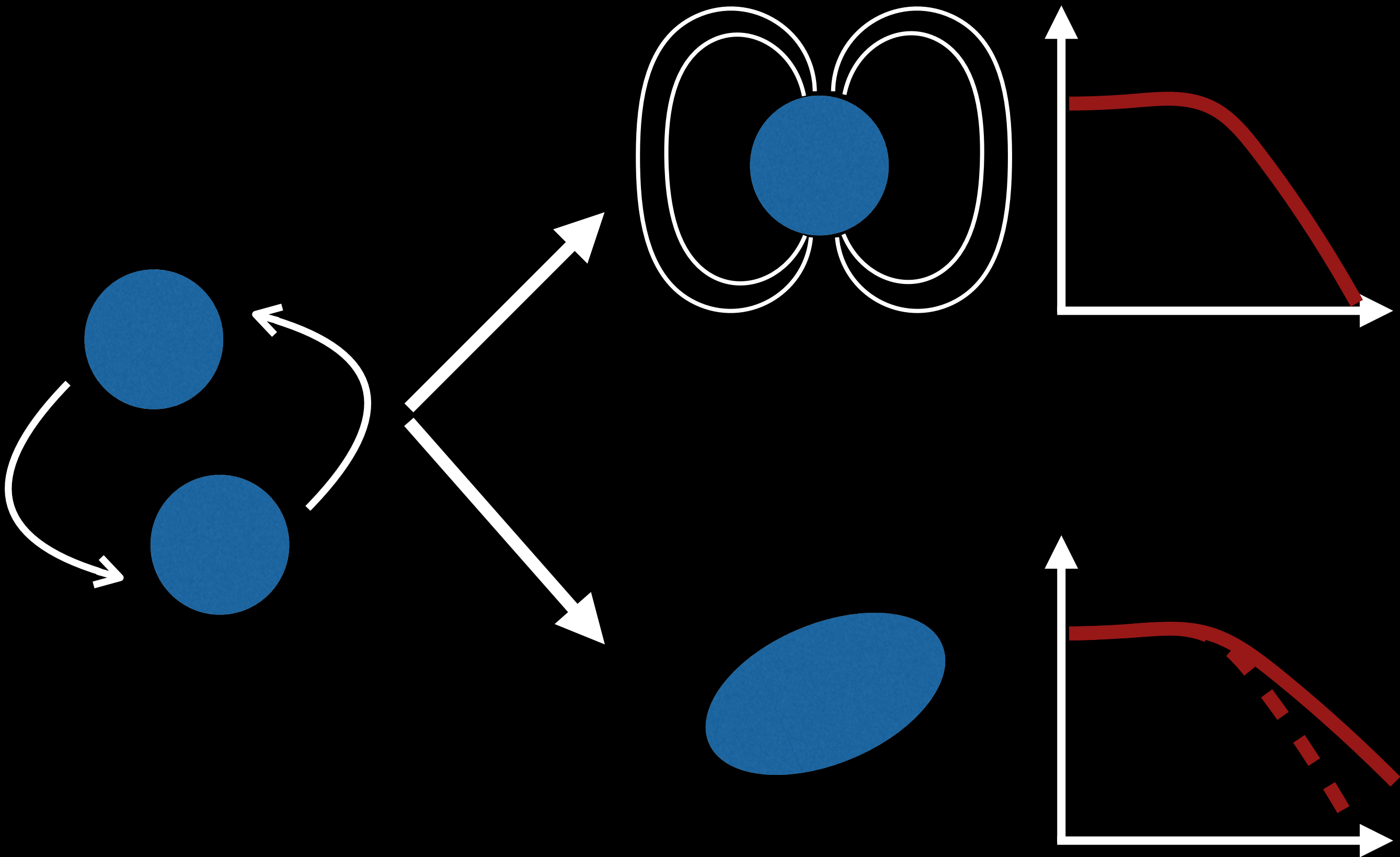


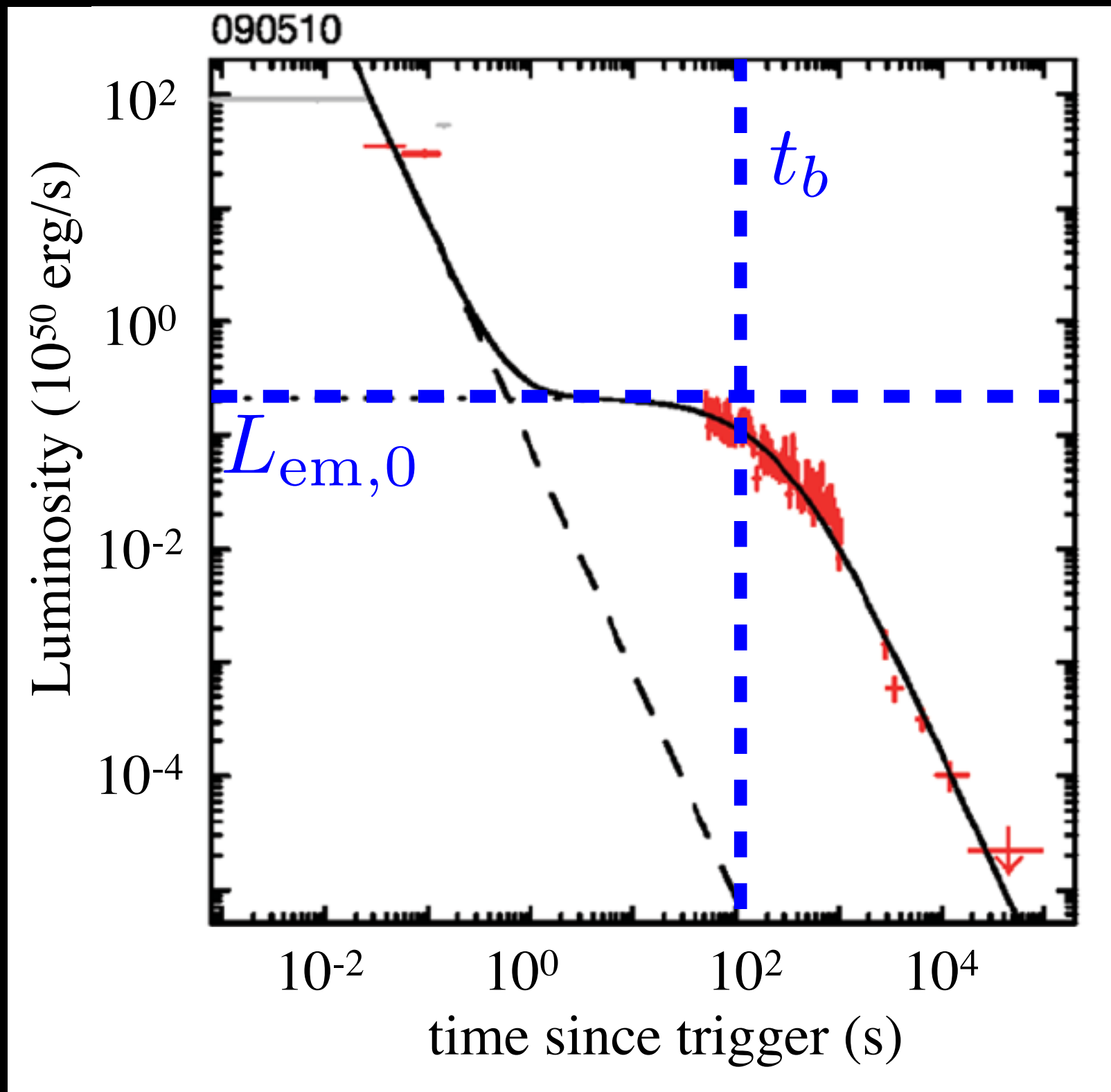
$\approx 1 \text{ yr}^{-1}$

ISS-Lobster

**What else can
we learn?**

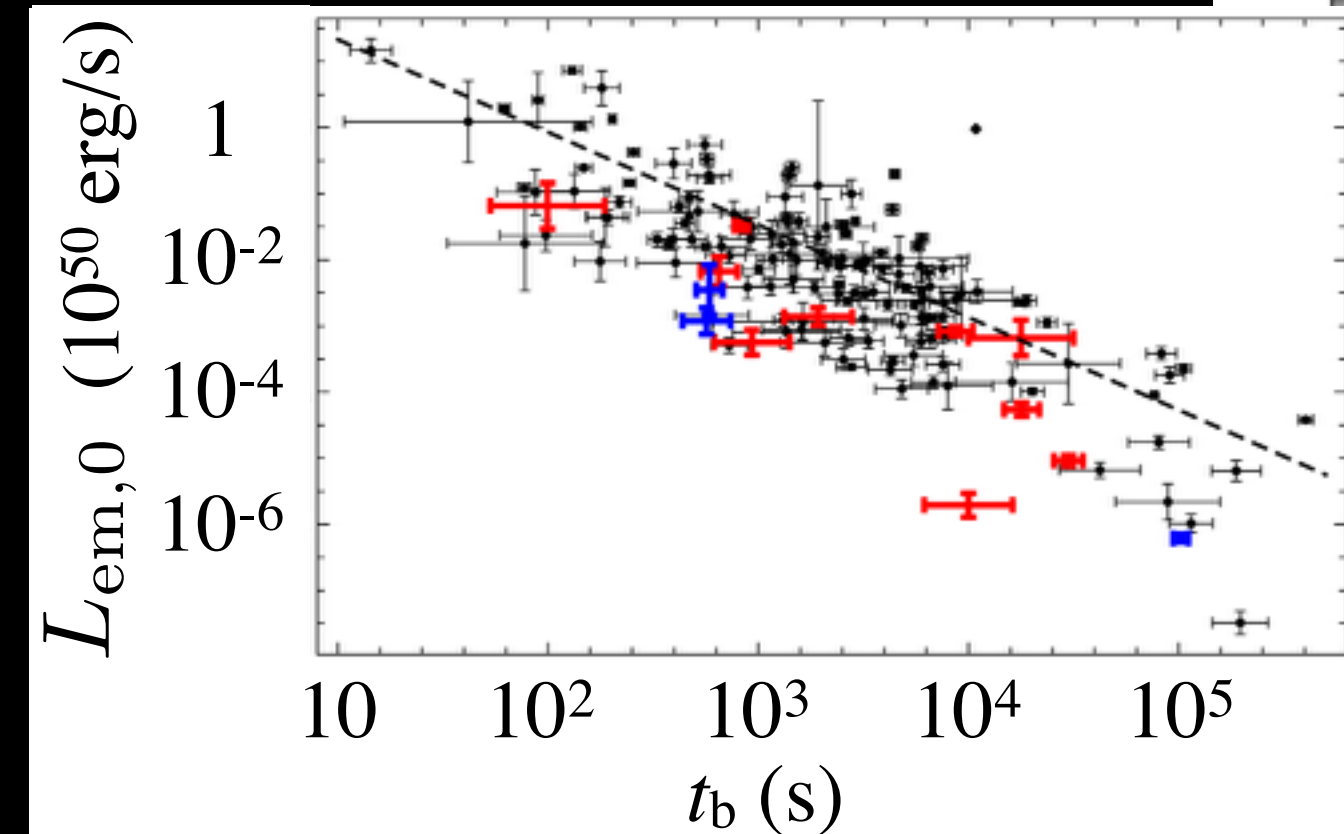
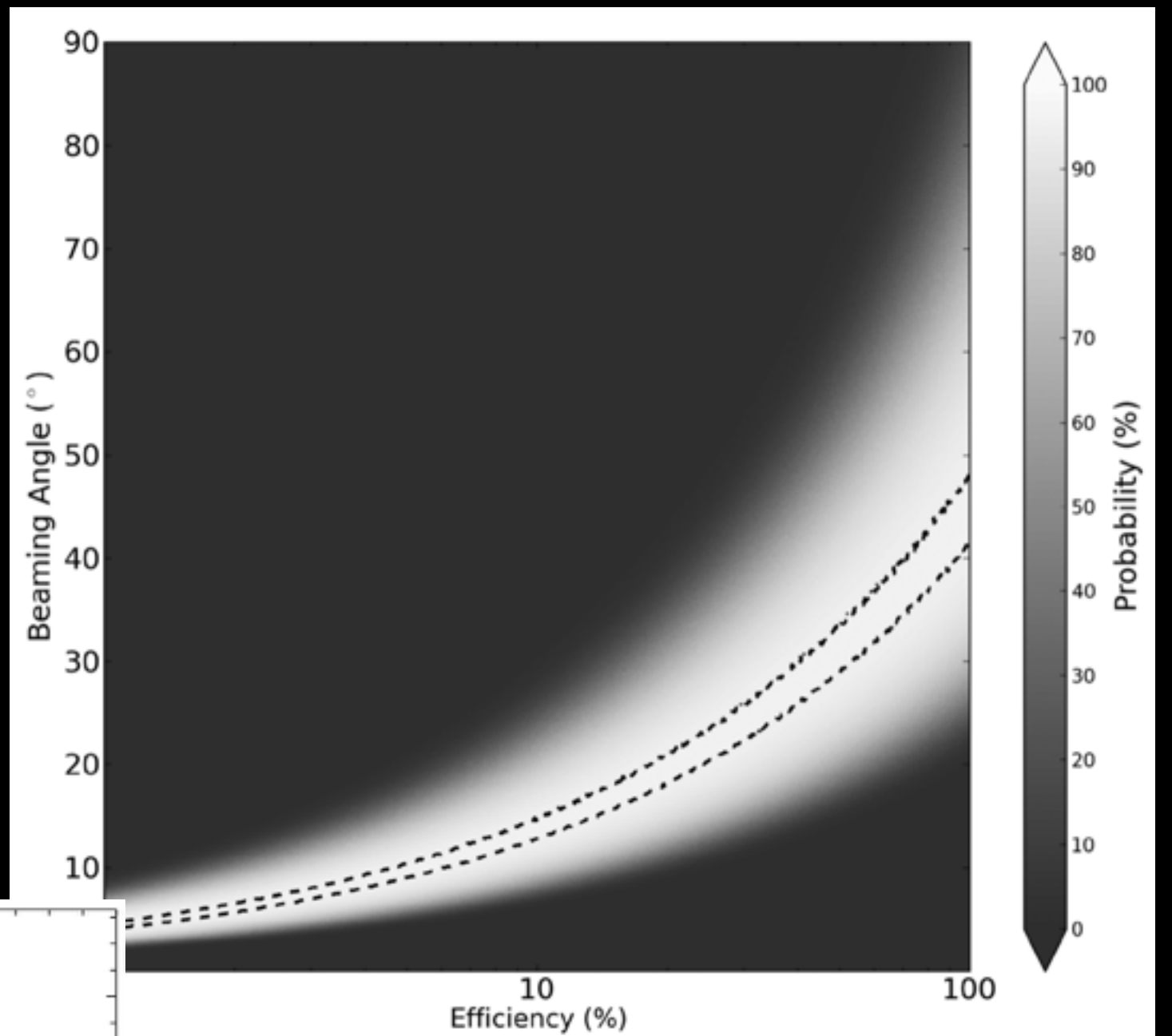






$$\epsilon \leq 0.33\eta \left(\frac{I}{10^{45} \text{ g cm}^2} \right)^{1/2} \left(\frac{L_{\text{em},0}}{10^{49} \text{ erg s}^{-1}} \right)^{-1} \left(\frac{t_b}{100 \text{ s}} \right)^{-3/2}$$

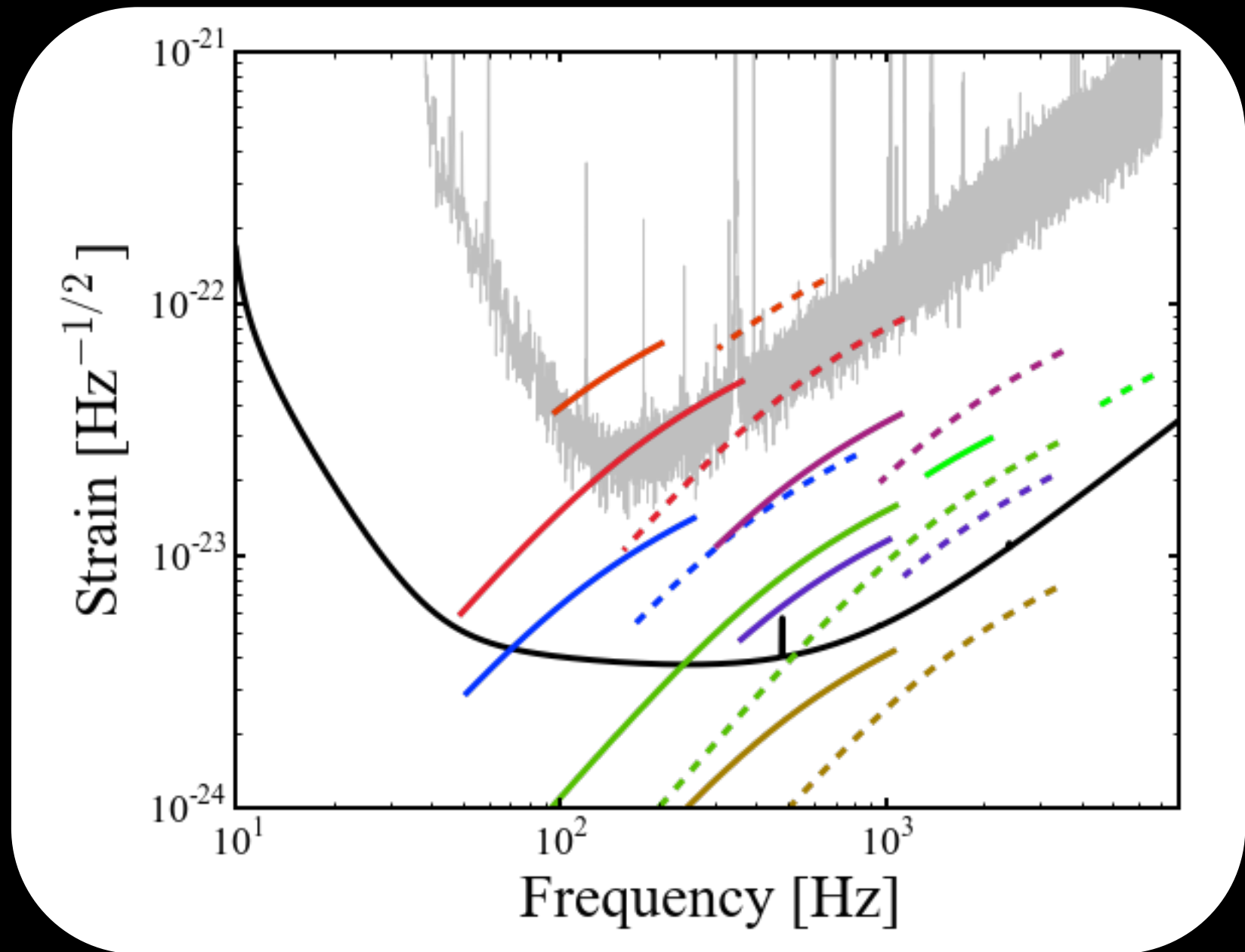
$$L_{\text{em},0} \propto \frac{\eta}{1 - \cos \theta} \frac{1}{t_b}$$



Rowlinson et al. (2014)

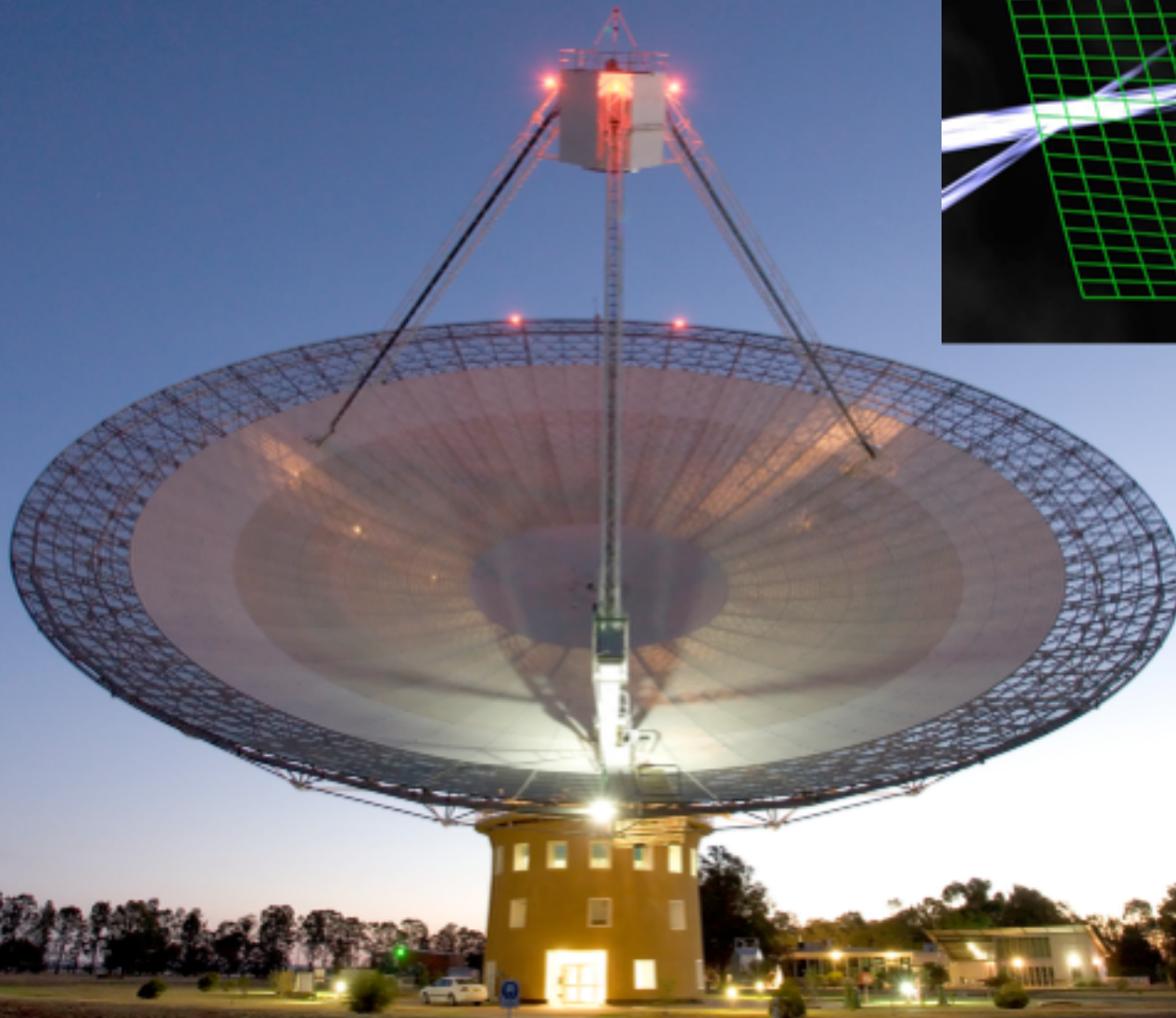
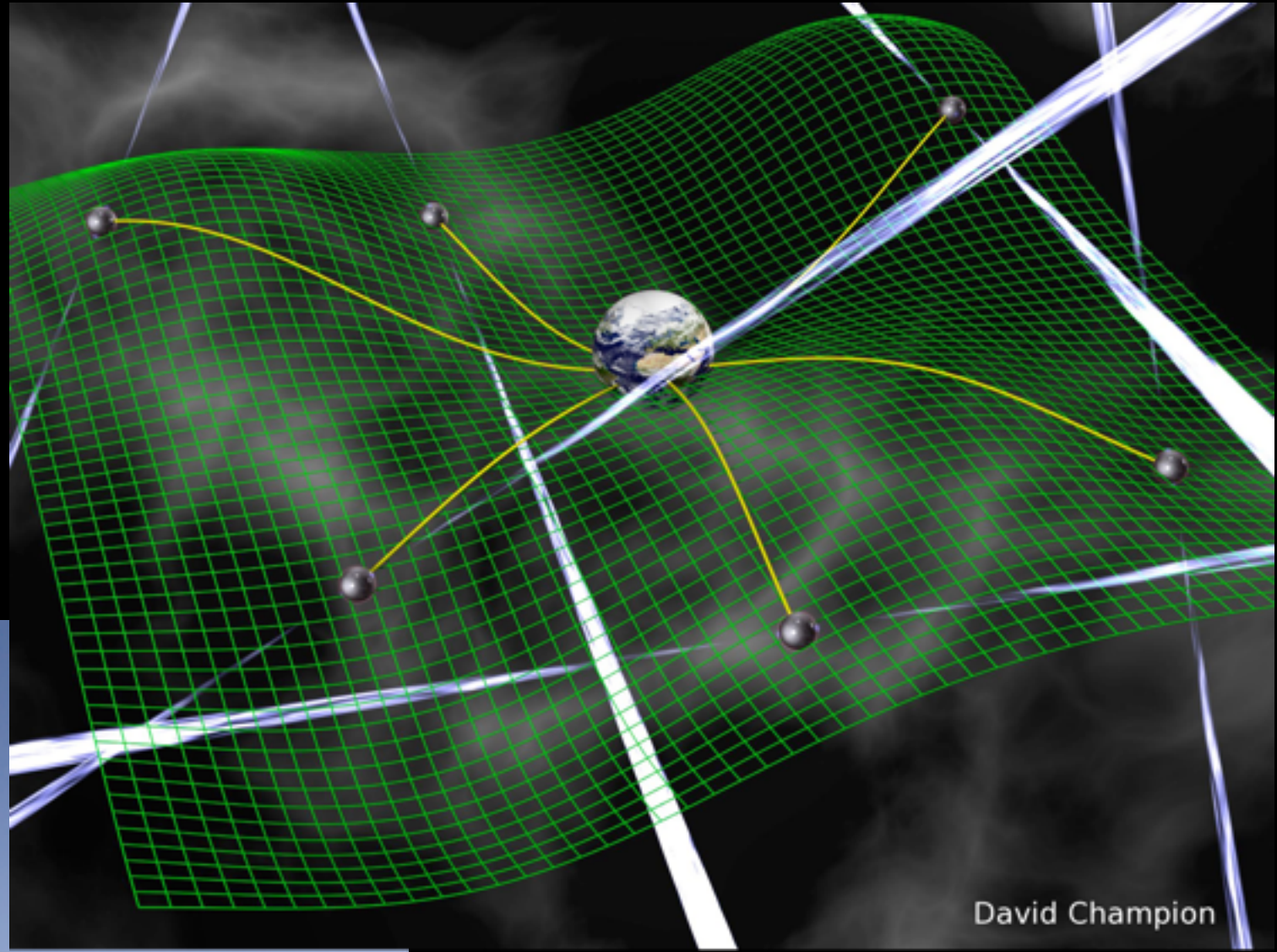
efficiency: $\eta = 0.1$

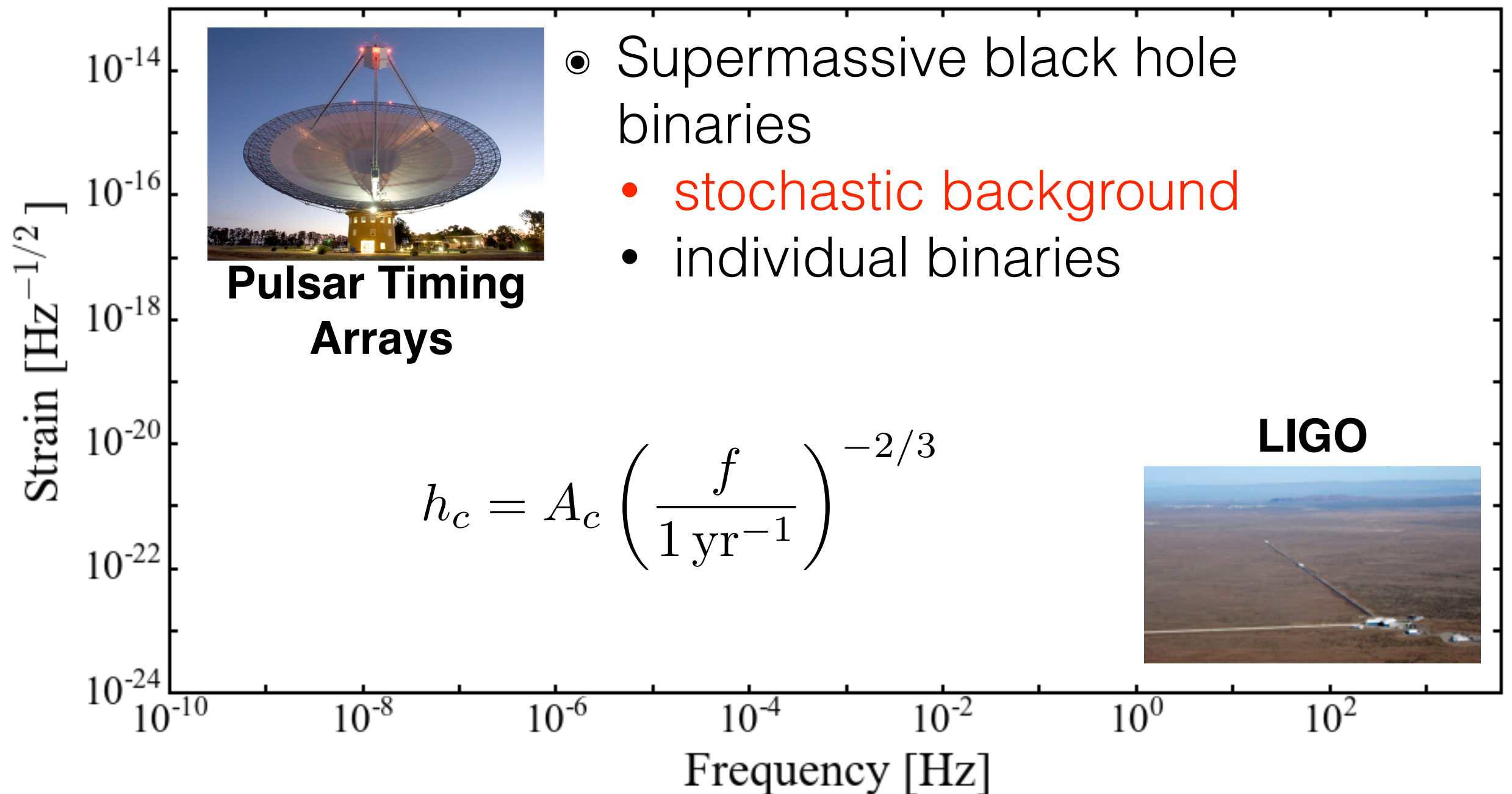
GRB	ϵ
GRB051221A	0.0061
GRB060801	0.0094
GRB070724A	0.0223
GRB070809	0.0049
GRB080905A	0.1532
GRB090426	0.0039
GRB090510	0.0041
GRB101219A	0.0012



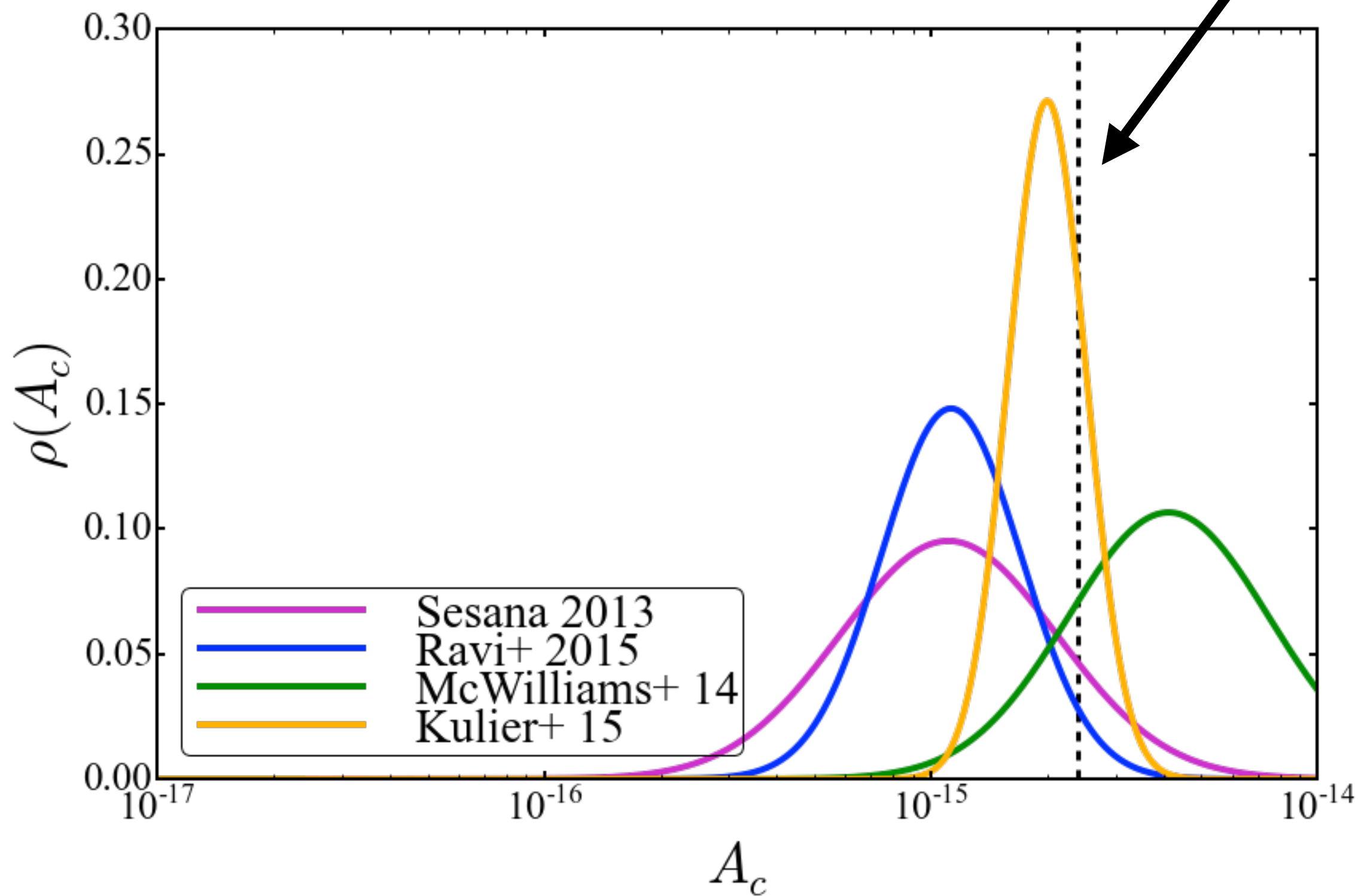


Parkes Pulsar Timing Arrays



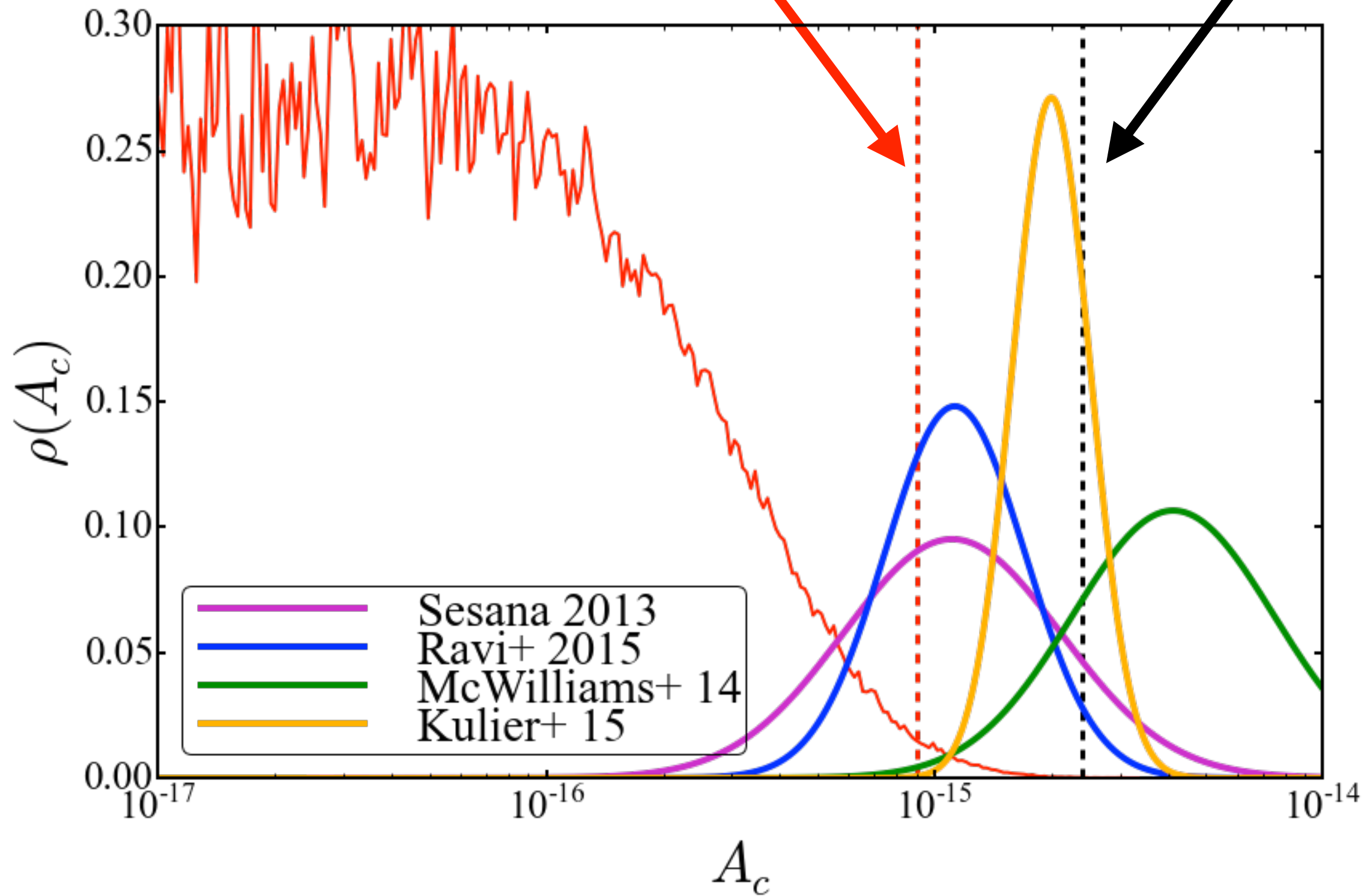


**Shannon et al.
(PPTA; 2013)**



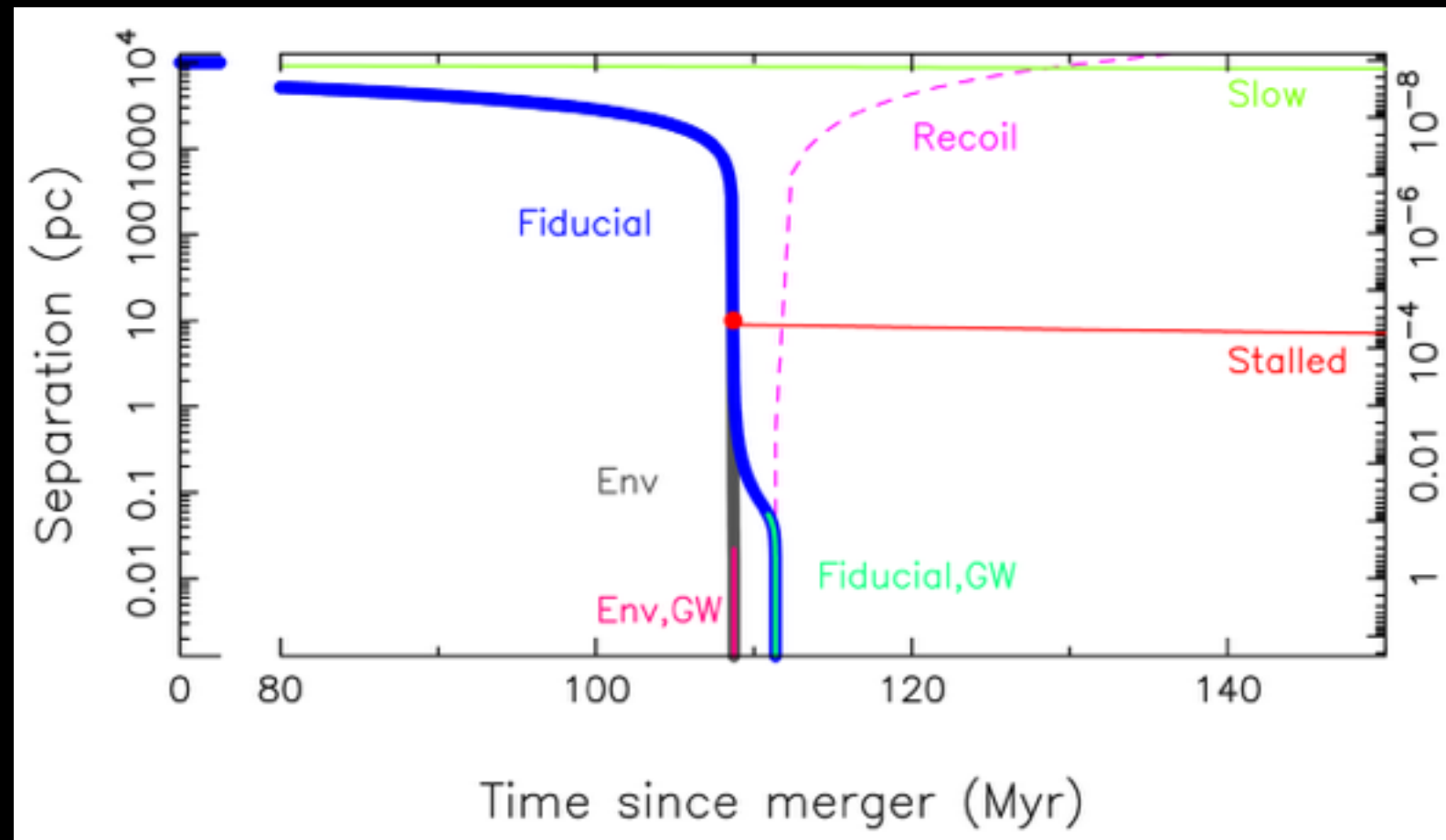
**Shannon et al.
(PPTA; 2015)**

**Shannon et al.
(PPTA; 2013)**



Astrophysical Inference

- Massive end of galaxy mass function?
- Black hole — bulge relations?
- Galaxy merger rate?
- Environmental factors: stars, gas, ... ?



Shannon et al.
(PPTA; submitted)

Conclusions

- If **millisecond magnetar** model is correct
(doesn't work with Rezzolla & Kumar (2015), Ciolfi & Siegel (2015) model)
 - allows for **equation of state** measurements
 - gravitational wave observations will help
 - allows us to study (in a limited sense) neutron star dynamics
- Parkes Pulsar Timing Array at 'design sensitivity'
 - no detection; important cosmological implications