

Flory-Huggins free energy (per site)  $f$

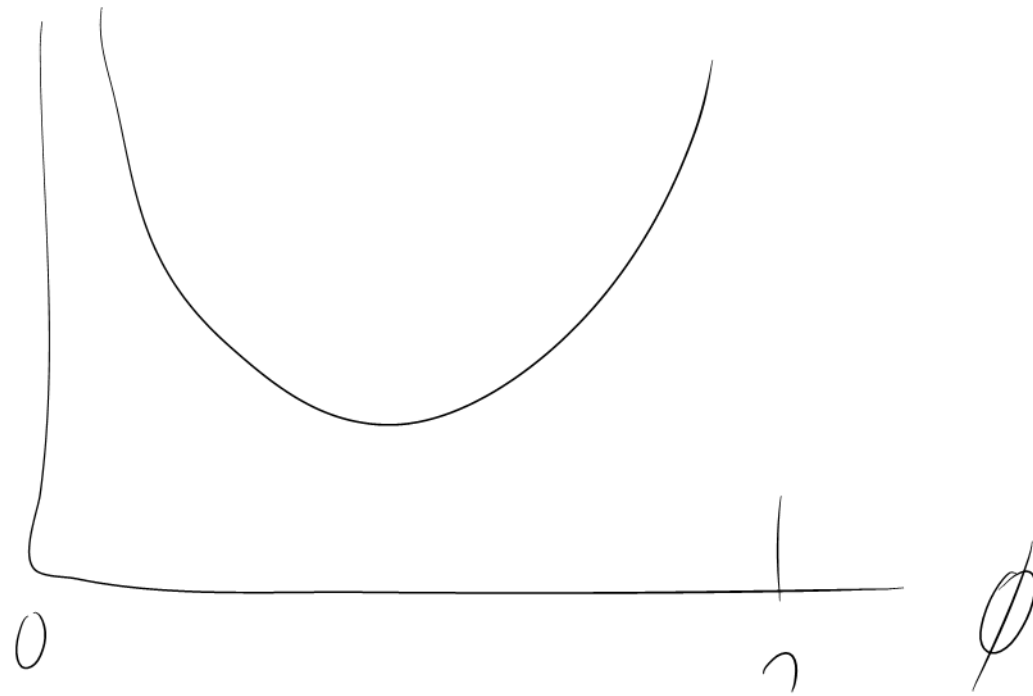
$$\frac{f}{k_B T} = \frac{\phi_A}{N_A} \ln \frac{\phi_A}{N_A} + \frac{\phi_B}{N_B} \ln \frac{\phi_B}{N_B} + \chi \phi_A \phi_B$$

+ terms linear in  $\phi$  + const.

do not matter for the  
phase behavior

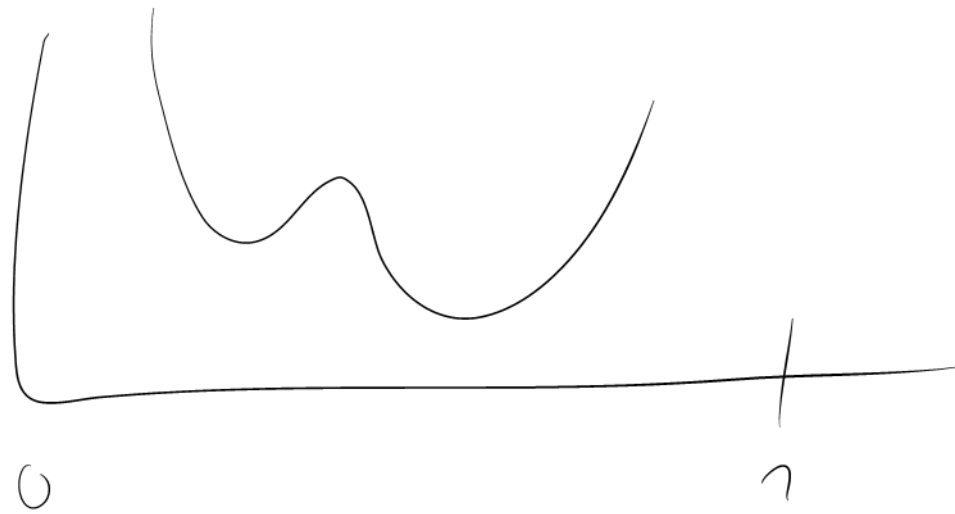
$$\phi = \phi_A = 1 - \phi_B$$

$f(\phi)$



$\lambda$  small

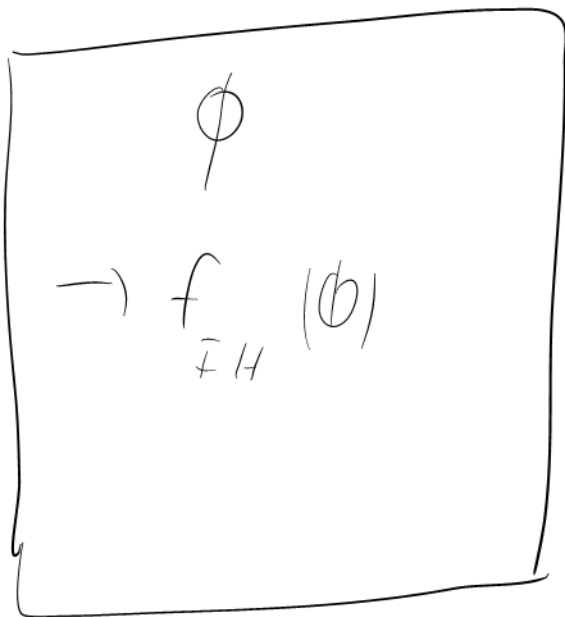
$f(\phi)$



$\lambda$  large

→ phase  
separation

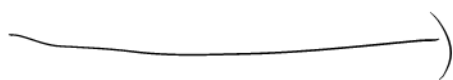
phase separation:



$f_{FH}(\lambda\phi_I + (1-\lambda)\phi_{II})$   
is the

$\phi$  conservation  
"lever rule"

$$\phi = \lambda\phi_I + (1-\lambda)\phi_{II}$$

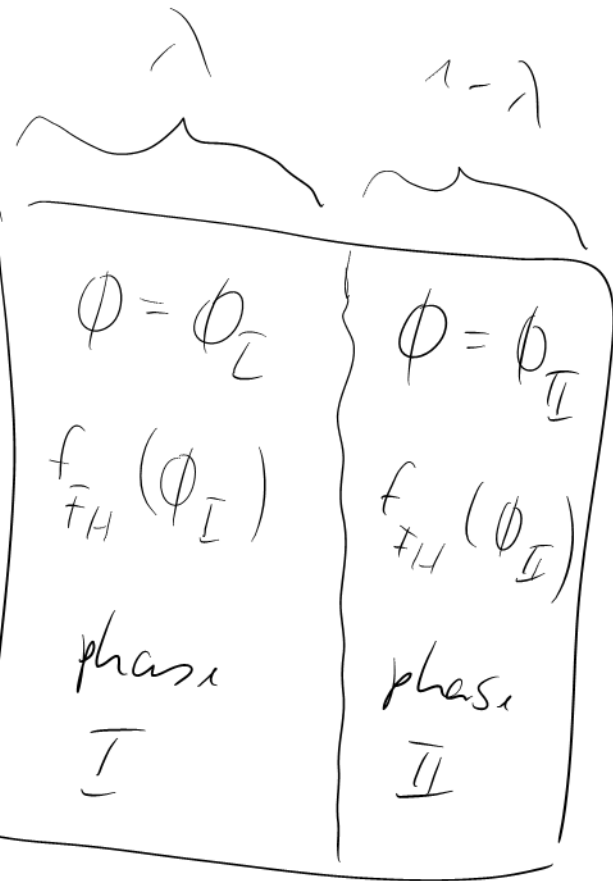


$\phi_I, \phi_{II}$   
determined  
by thermo-  
dynamics



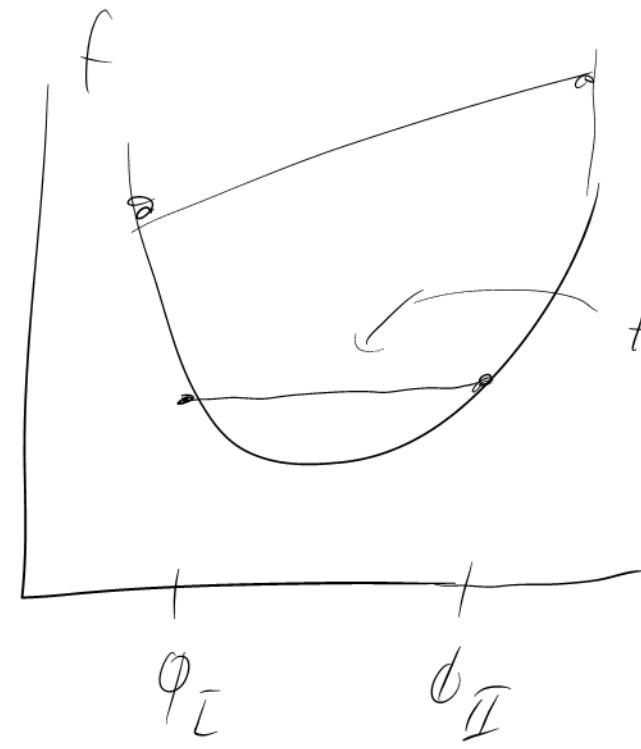
condition for phase separation

$$\lambda \in [0, 1]$$



$$\lambda f_{FH}(\phi_I) + (1-\lambda) f_{FH}(\phi_{II})$$

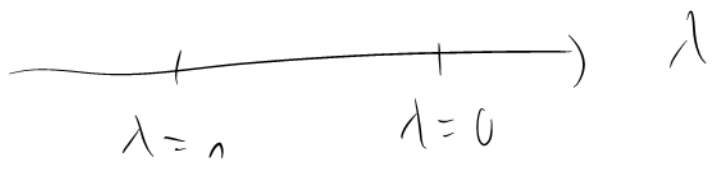
f convex:



$$\phi = \lambda \phi_I + (1-\lambda) \phi_{II}$$

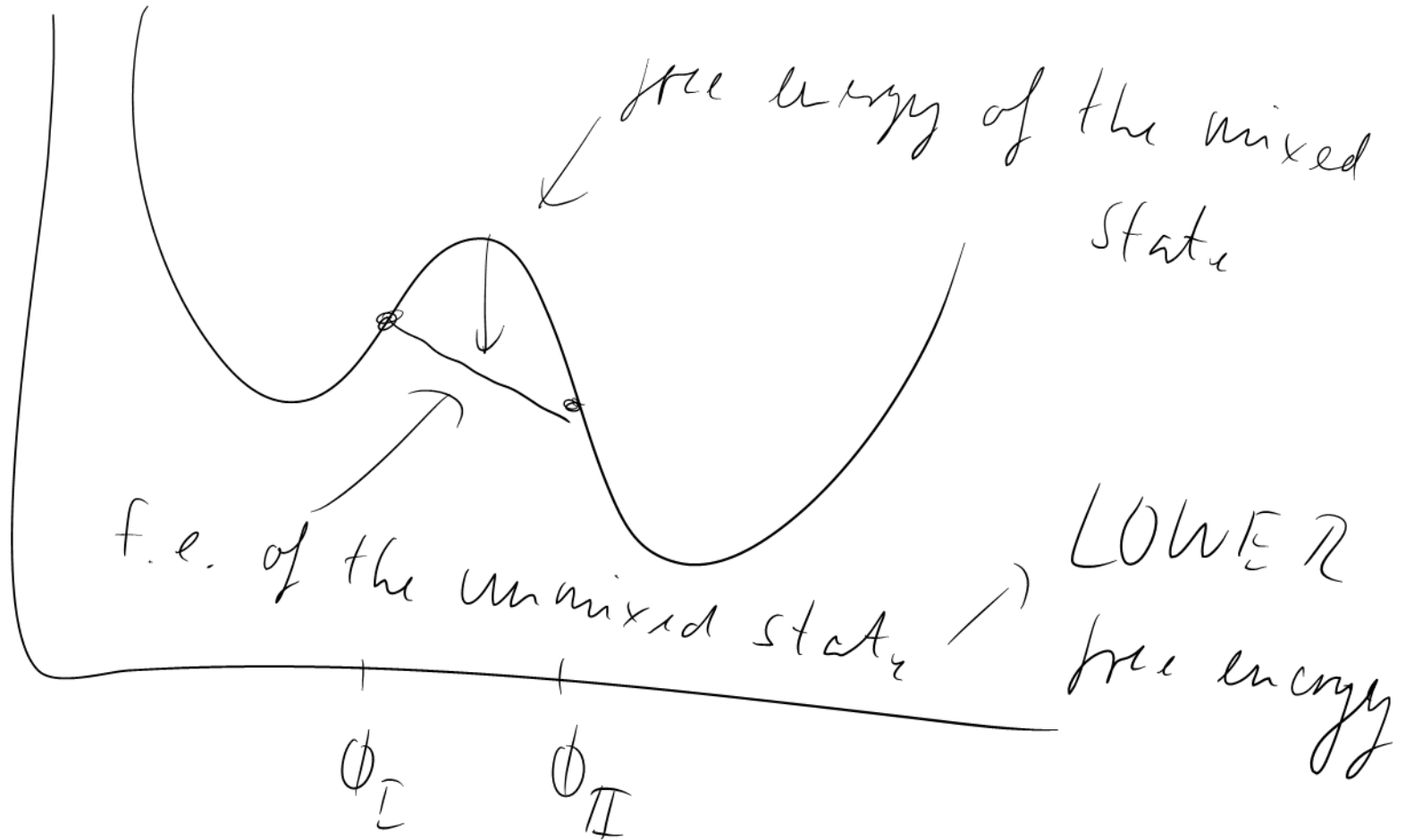
$$f(\phi_I) \cdot \lambda + f(\phi_{II}) \cdot (1-\lambda)$$

$\phi$  free energy of the unmixed state is larger

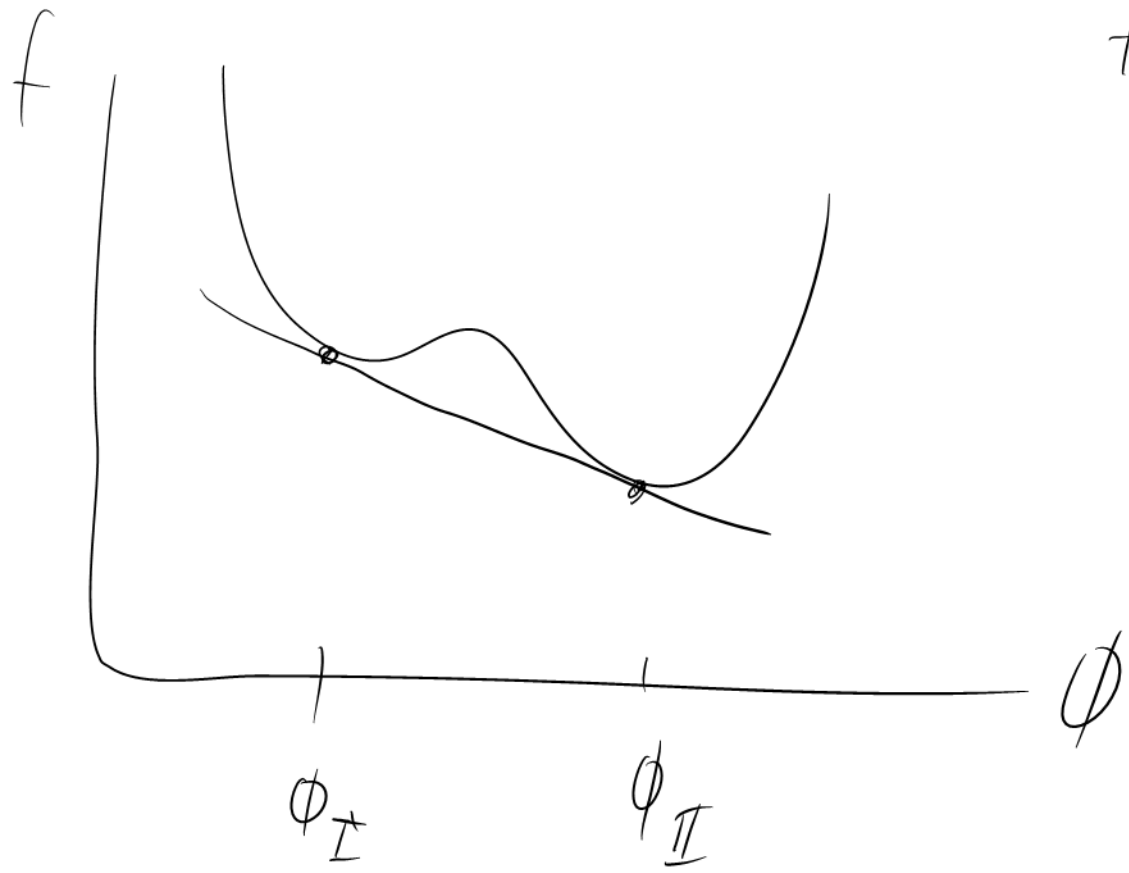


→ homogeneous state is stable!

f concave



optimize  $\phi_I, \phi_{II} \rightarrow$  common target



this determines  
 $\phi_I, \phi_{II}!$

$$\frac{\partial f}{\partial \phi} \Big|_{\phi = \phi_I} = \frac{\partial f}{\partial \phi} \Big|_{\phi = \phi_{II}} = \frac{f(\phi_I) - f(\phi_{II})}{\phi_I - \phi_{II}}$$

2 eq., 2 unknowns ( $\phi_I, \phi_{II}$ )

$$f \rightarrow f + A\phi + B$$

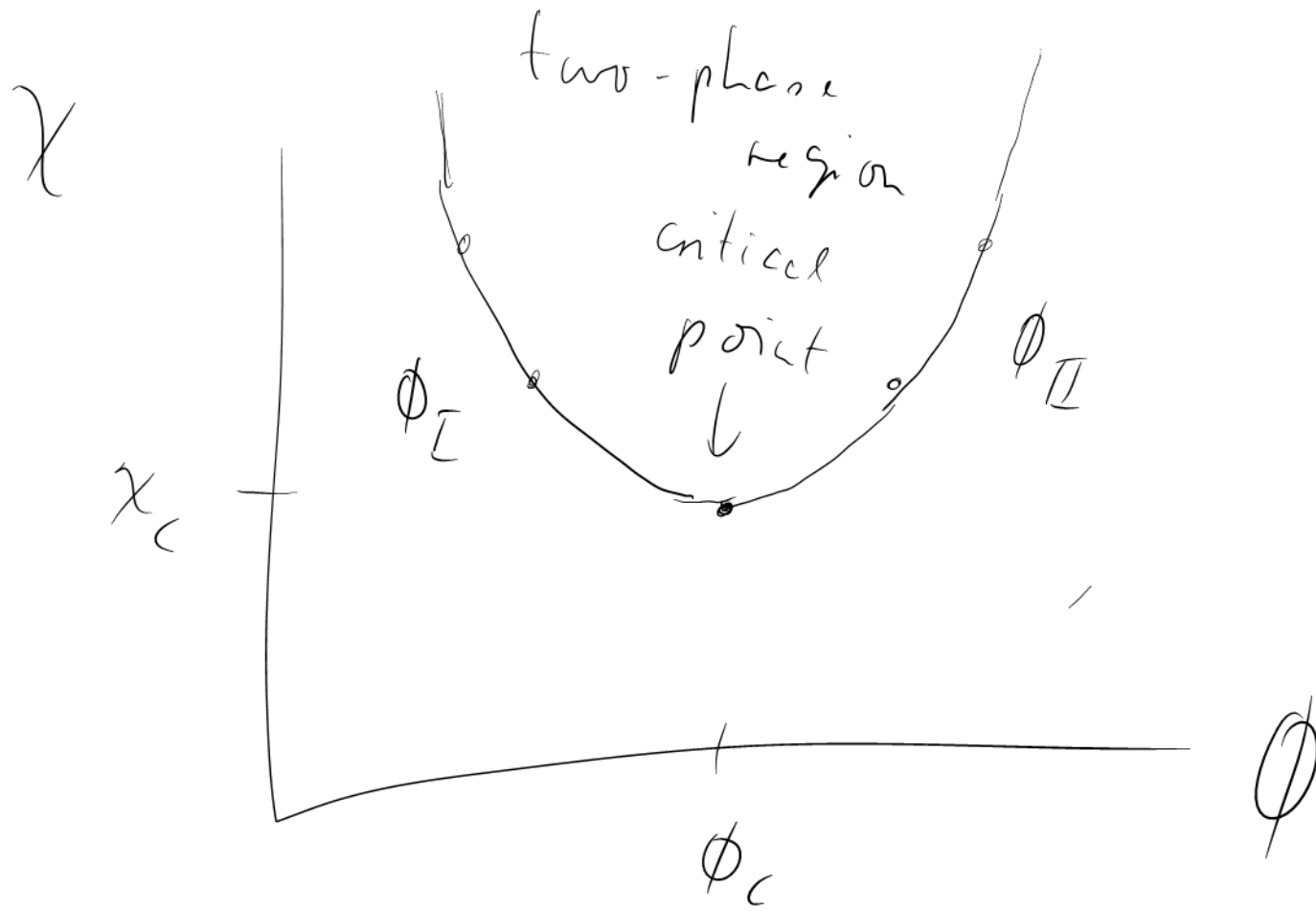
$$\frac{\partial f}{\partial \phi} \Big|_{\phi_I} \rightarrow \frac{\partial f}{\partial \phi} \Big|_{\phi_I} + A \quad \frac{\partial f}{\partial \phi} \Big|_{\phi_{II}} \rightarrow \frac{\partial f}{\partial \phi} \Big|_{\phi_{II}} + A$$

$$f(\phi_I) - f(\phi_{II}) \rightarrow f(\phi_{II}) - f(\phi_{II}) + A(\phi_I - \phi_{II})$$

$$\frac{f(\phi_I) - f(\phi_{II})}{\phi_I - \phi_{II}} \rightarrow \frac{f(\phi_I) - f(\phi_{II})}{\phi_I - \phi_{II}} + A$$

therefore we can neglect  $A\phi + B$  in  $f$ !





phase  
diagram

$$\frac{\partial^2 f}{\partial \phi^2} > 0$$

locally convex

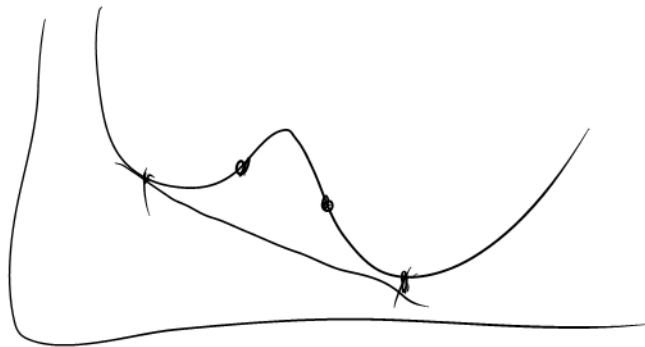
→ system is stable w.r.t.

infinitesimal fluctuations

$$\frac{\partial^2 f}{\partial \phi^2} < 0$$

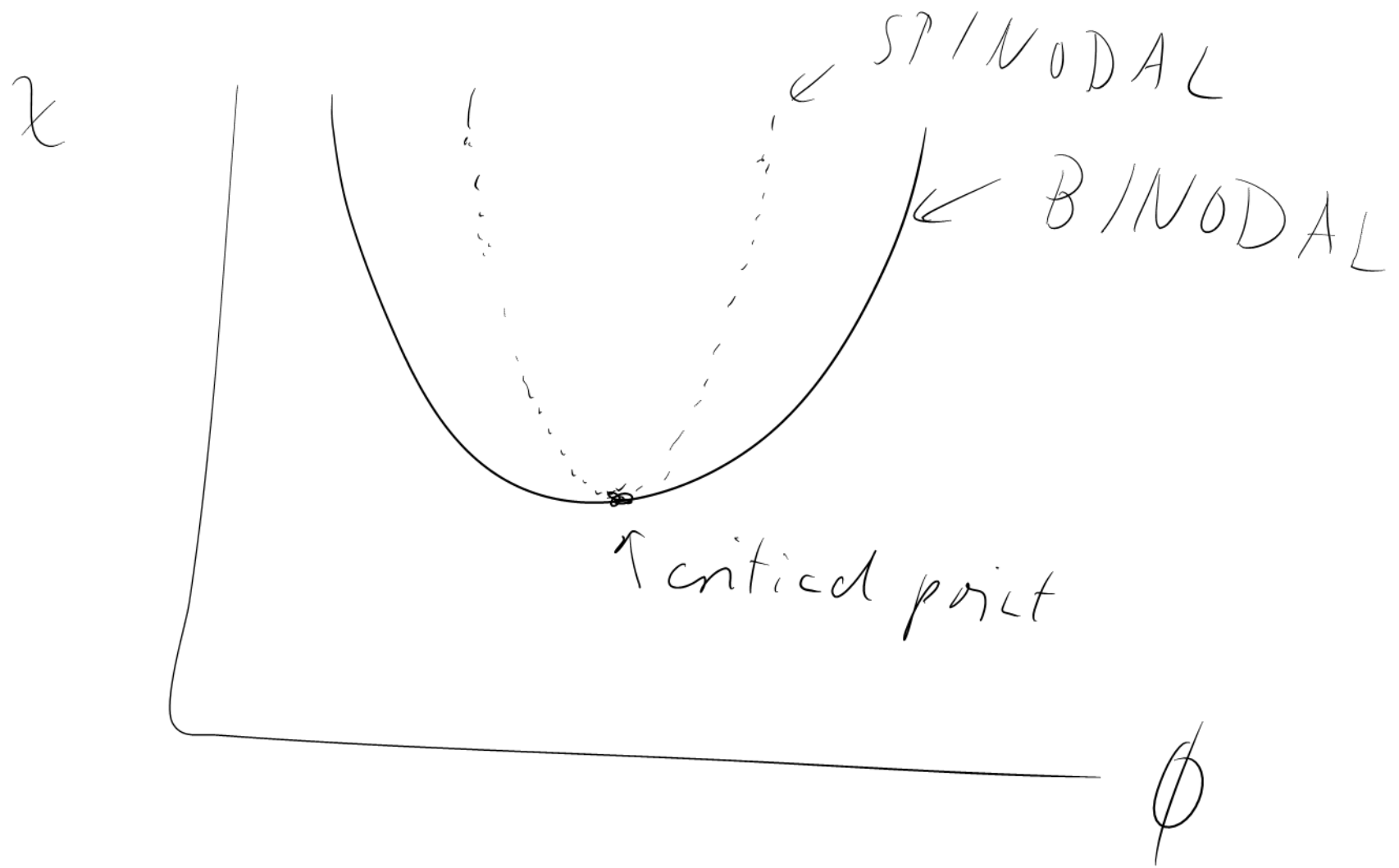
locally concave → unstable

even w.r.t. infinitesimal fluctuations



$$\frac{\partial^2 f}{\partial \phi^2} = 0 \quad \text{inflection points}$$

→ SPINODAL



(phase diagram again)

# Nucleation

(excursion)

# Burke-Döring



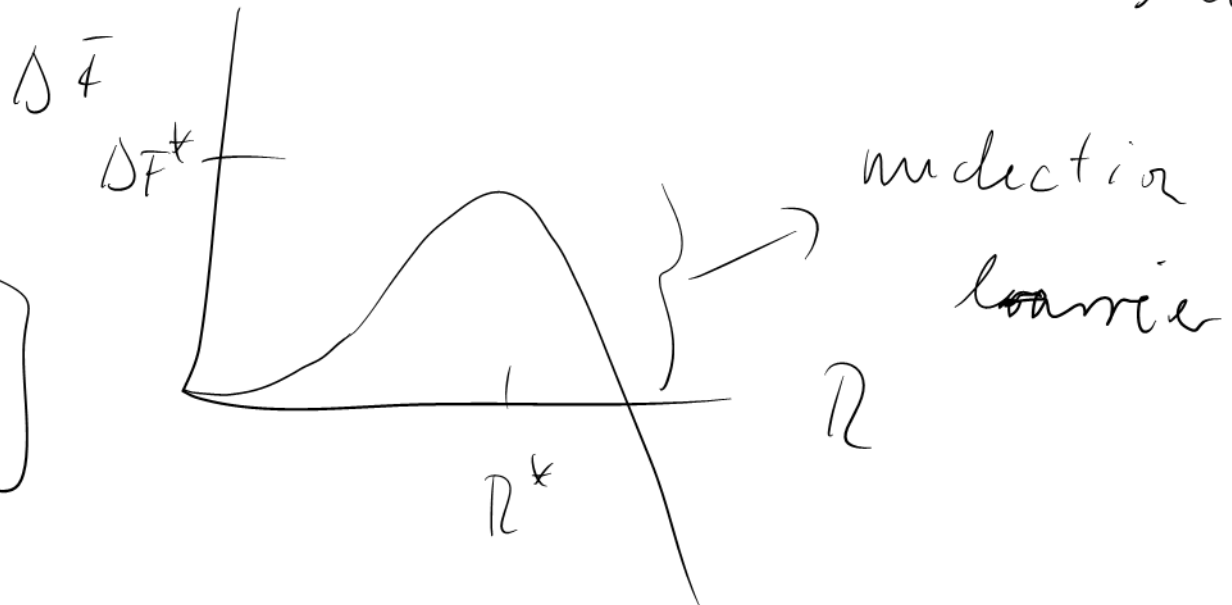
Phase II

$$\Delta F = -\frac{4\pi}{3} R^3 \cdot \Delta\mu$$

$$+ 4\pi R^2 \cdot \gamma \leftarrow \text{interfacial tension}$$

time ~

$$\tau \propto \exp\left[ + \frac{\Delta F^*}{k_B T} \right]$$

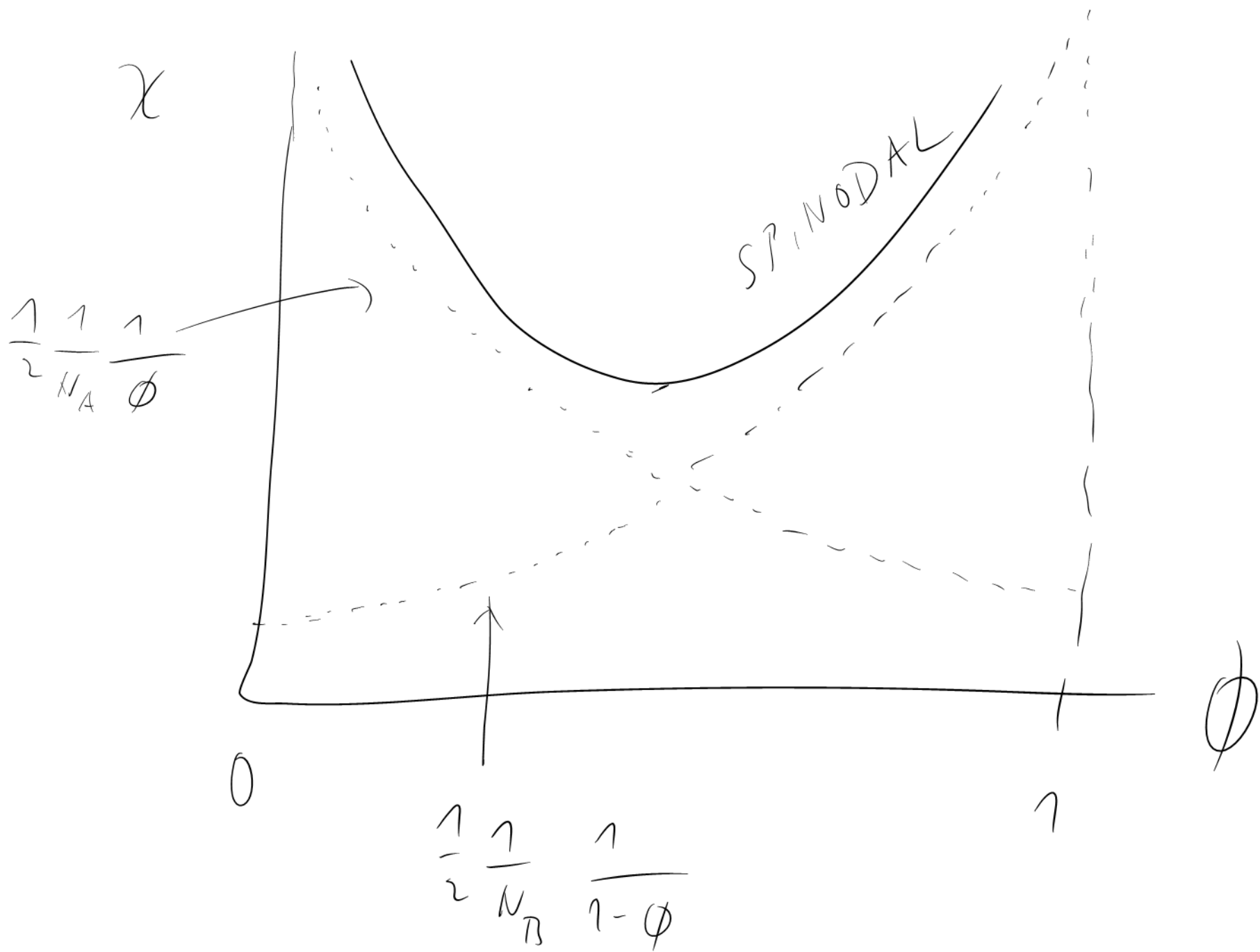


back to FH: calculated spinodal

$$\frac{f}{k_B T} = \frac{\phi}{N_A} \ln \frac{\phi}{N_A} + \frac{1-\phi}{N_B} \ln \frac{1-\phi}{N_B} + \chi \phi (1-\phi)$$

$$\frac{\partial^2}{\partial \phi^2} \frac{f}{k_B T} = \frac{1}{N_A \phi} + \frac{1}{N_B (1-\phi)} - 2\chi \stackrel{!}{=} 0$$

$$\chi = \frac{1}{2} \left( \frac{1}{N_A \phi} + \frac{1}{N_B (1-\phi)} \right)$$



critical point?

$$0 \stackrel{!}{=} \frac{\partial}{\partial \phi} \left\{ \frac{1}{2} \left( \frac{1}{N_A \phi} + \frac{1}{N_B (1-\phi)} \right) \right\}$$

$$= \frac{\partial}{\partial \phi} \left\{ \frac{1}{N_A \phi} + \frac{1}{N_B (1-\phi)} \right\} = -\frac{1}{N_A \phi^2} + \frac{1}{N_B (1-\phi)^2}$$

$$\Rightarrow \phi = \phi_c = \frac{\sqrt{N_B}}{\sqrt{N_A} + \sqrt{N_B}} \quad \rightarrow \text{into spinodal}$$

$$\Rightarrow \chi_c = \frac{1}{2} \left( \frac{1}{N_A} + \frac{1}{N_B} \right) + \frac{1}{\sqrt{N_A N_B}}$$

Symmetric polymer mixture:  $N_A = N_B = N$

$$x_c = \frac{1}{N} + \frac{1}{N} = \frac{2}{N} \quad \phi_c = \frac{1}{2}$$

↓  
 $\propto \frac{1}{k_B T_c}$

→  $k_B T_c \propto N$

polymers  
mix very  
poorly!



polymer solution

$$N_B = 1$$

$$N_A \gg 1$$

$$N_A = N$$

$$\phi_c \approx \frac{1}{\sqrt{N}} = N^{-1/2}$$

$$\chi_c \approx \frac{1}{2} + \frac{1}{\sqrt{N}} = \frac{1}{2} + N^{-1/2}$$