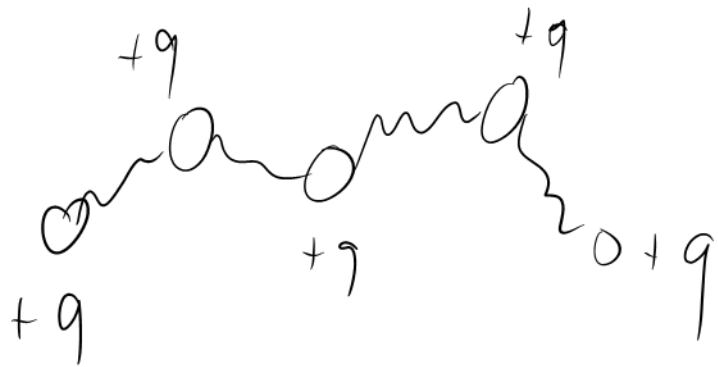


9. Polyelectrolyte Conformations I:

Gaussian Chain + Charges



ions and other polymers
infinitely far away

$$\mathcal{H} = \underbrace{\frac{3k_B T}{2b^2} \sum_{i=1}^{N-1} (\vec{r}_{i+1} - \vec{r}_i)^2}_{\text{bead-spring model}} + \frac{q^2}{4\pi\epsilon} \sum_{i < j} \frac{1}{|\vec{r}_i - \vec{r}_j|}$$

Coulomb repulsion

Set $\eta = fe$ f : "fraction"

Weak damping: $f^2 u \ll 1$

$u = \frac{l_B}{b}$ coupling parameter

$$\beta \mathcal{H} = \frac{3}{2b^2} \sum_{i=1}^{N_s} (\vec{r}_{i+1} - \vec{r}_i)^2 + f^2 l_B \sum_{i < j} \frac{1}{|\vec{r}_i - \vec{r}_j|}$$

uncharged

$$\vec{R}_{FE} \equiv \vec{r}_N - \vec{r}_1$$

$$\langle \vec{R}_{FE}^2 \rangle = b^2 (N-1) \approx b^2 N$$

charged: chain strongly stretched in some direction, due to electrostatic repulsion

intermission:



→ "Pinches blobs"
stretching by an external force

$$\mathcal{H} = \frac{3k_B T}{2b^2} \sum_{i=1}^{N-1} (\vec{r}_{i+1} - \vec{r}_i)^2 - \underbrace{\vec{F} \cdot \vec{r}_N + \vec{F} \cdot \vec{r}_1}_{-\vec{F} \cdot (\vec{r}_N - \vec{r}_1)}$$

$$\begin{aligned} \vec{r}_N - \vec{r}_1 &= (\vec{r}_N - \vec{r}_{N-1}) + (\vec{r}_{N-1} - \vec{r}_{N-2}) + \dots + (\vec{r}_2 - \vec{r}_1) \\ &= \sum_{i=1}^{N-1} (\vec{r}_{i+1} - \vec{r}_i) \end{aligned}$$

$$\mathcal{H} = \frac{3k_B T}{2b^2} \sum_{i=1}^{N-1} \left\{ (\vec{r}_{i+1} - \vec{r}_i)^2 - \frac{2b^2}{3k_B T} \vec{F} \cdot (\vec{r}_{i+1} - \vec{r}_i) \right\} =$$

$$= \frac{3k_B T}{2b^2} \sum_{i=n}^{N-n} \left\{ \left(\vec{r}_{i+n} - \vec{r}_i - \frac{b^2 \vec{F}}{3k_B T} \right)^2 - \left(\frac{b^2 \vec{F}}{3k_B T} \right)^2 \right\}$$

= const., discard

$$\rightarrow \frac{3k_B T}{2b^2} \sum_{i=n}^{N-n} \left(\vec{r}_{i+n} - \vec{r}_i - \frac{b^2 \vec{F}}{3k_B T} \right)^2$$

$$\langle \vec{r}_{i+n} - \vec{r}_i \rangle = \frac{b^2 \vec{F}}{3k_B T} ; \text{ equipartition theorem } \Rightarrow$$

$$\frac{3k_B T}{2b^2} \left\langle \left(\vec{r}_{i+n} - \vec{r}_i - \frac{b^2 \vec{F}}{3k_B T} \right)^2 \right\rangle = \frac{3}{2} k_B T$$

$$\left\langle \left(\vec{r}_{i+n} - \vec{r}_i - \frac{b^2 \vec{F}}{3k_B T} \right)^2 \right\rangle = b^2$$

$$\overline{r_{i+1}} - \overline{r_i} = \frac{b^2 \vec{f}}{3 k_B T} + \frac{1}{\sqrt{3}} b \vec{q}_i$$

↑ Gaussian random vector

$$\langle \vec{q}_i \rangle = 0$$

$$\langle q_{i\alpha} q_{j\beta} \rangle = \delta_{ij} \delta_{\alpha\beta}$$

$$\overline{R_{\overline{10}}} = \overline{r_N} - \overline{r_0} = \sum_{i=0}^{N-1} (\overline{r_{i+1}} - \overline{r_i}) =$$

$$= \underbrace{(N-1) \frac{b^2 \vec{f}}{3 k_B T}}_{O(N)} + \frac{b}{\sqrt{3}} \sum_{i=0}^{N-1} \vec{q}_i \rightarrow O(N^{1/2})$$

extension of chain

$$R_{E||} \sim N \frac{b^2 F}{3k_B T}$$

$$R_{E\perp} \sim \sqrt{\frac{2}{3}} N^{1/2} b$$

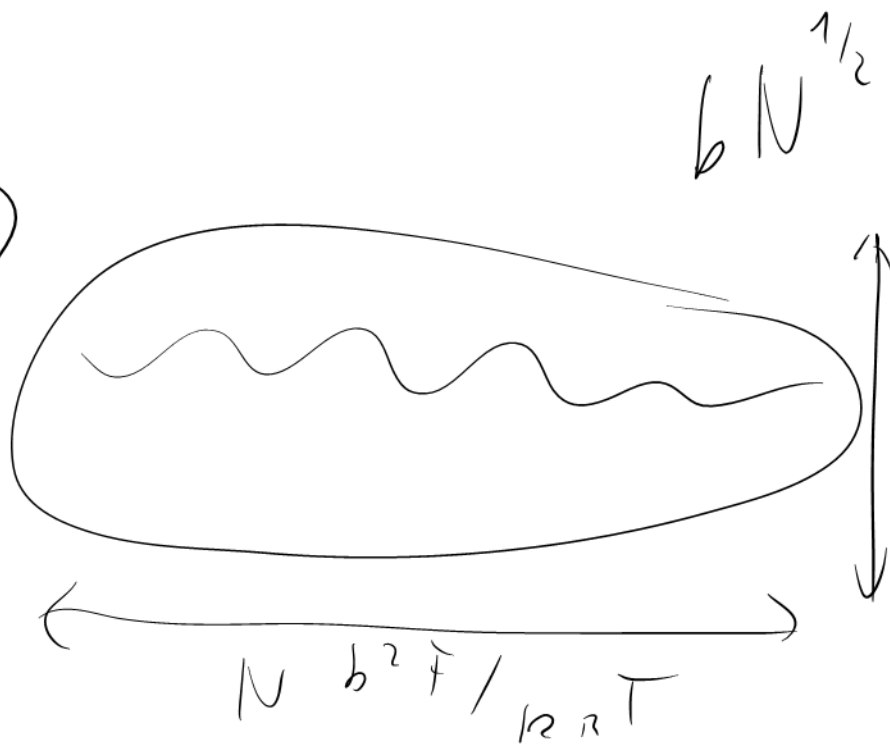
$$\sqrt{\langle R_{E||}^2 \rangle}$$

$$\sqrt{\langle R_{E\perp}^2 \rangle}$$

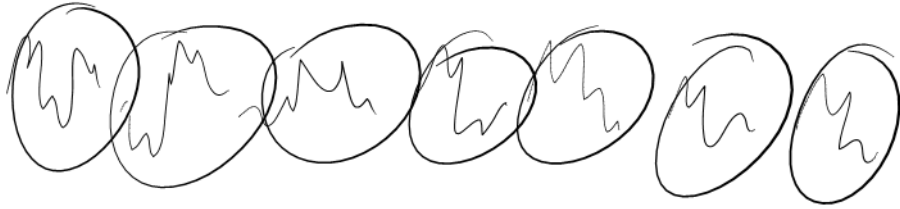
$$b^2 = \langle (\vec{r}_{i+1} - \vec{r}_i)^2 \rangle$$

$$\uparrow$$

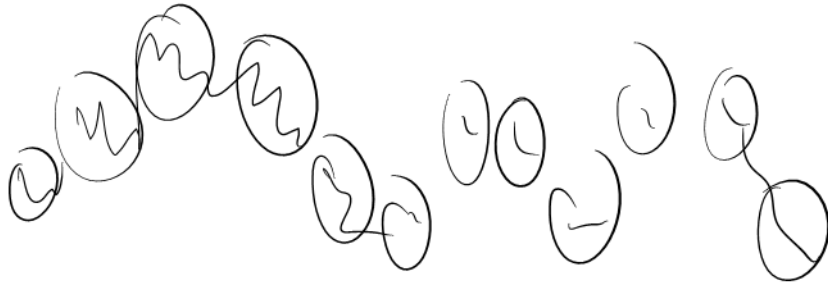
$$F_{||} = 0$$



"Pinch blobs"



\mathbb{R}^2



one blob: n monomers

$$\# \text{ blobs: } \frac{N}{n}$$

within a blob: $RW \Rightarrow \xi \sim b n^{1/2}$

$$F \xi \sim k_B T \quad \xi \sim \frac{k_B T}{F} \quad b n^{1/2} \sim \frac{k_B T}{F}$$

$$n \sim \left(\frac{k_B T}{F b} \right)^2, \quad \text{number of blobs} \sim \frac{N}{n} \sim$$
$$\sim N \left(\frac{F b}{k_B T} \right)^2$$

$$R_{E||} \sim \frac{N}{n} \left\{ - N \left(\frac{Fb}{k_B T} \right)^2 \frac{k_B T}{F} \sim N \frac{b^2 F}{k_B T} \right.$$

end of intermission, bc due to the changed
 dirn.

assume a stretched conformation,

extension R_E in one direction,

" $bN^{1/2}$ perpendicular

} a speed ratio
 $\frac{R_E}{bN^{1/2}}$

for electrostatics: assume geometry of an ellipsoid

$$\downarrow \text{electrostatic energy} \sim \frac{(Nfe)^2}{4\pi\epsilon R_E} \ln\left(\frac{R_E}{bN^{1/2}}\right)$$

$$\text{entropic elastic energy} \sim k_B T \frac{R_E^2}{Nb^2}$$

"Flory theory"

$$k_B T \frac{R_E^2}{Nb^2} + \frac{(Nfe)^2}{4\pi\epsilon R_E} \ln\left(\frac{R_E}{bN^{1/2}}\right) \stackrel{!}{=} \text{Min. wrt } R_E$$

$$\frac{R_E^2}{Nb^2} + \frac{N^2 f^2 \lambda_B}{R_E} \ln \left(\frac{R_E}{bN^{1/2}} \right) \stackrel{!}{=} \text{Min.}$$

first ignore log correction \rightarrow

$$\frac{R_E^2}{Nb^2} + \frac{N^2 f^2 \lambda_B}{R_E} \stackrel{!}{=} \text{Min}$$

$$\frac{R_E}{Nb^2} - \frac{1}{R_E^2} N^2 f^2 \lambda_B = 0 \quad | \cdot R_E^2 \cdot Nb^2$$

$$R_E^3 = N^3 b^2 f^2 l_B$$

coupling parameter $u = \frac{l_B}{b}$

$$R_E^3 = N^3 b^3 f^2 u$$

$$R_E \sim N b (f^2 u)^{1/3}$$

result for
neglected log-
correction

now, take log-connection into account

$$\text{Write: } R_E = Nb (f^2 n)^{1/3} x \quad \leftarrow \text{determine } x$$

$$\frac{R_E^2}{Nb^2} \sim N (f^2 n)^{2/3} x^2$$

$$\frac{N^2 f^2 l_B}{R_E} \sim N \underbrace{f^2 \frac{l_B}{b}}_{f^2 n} (f^2 n)^{-1/3} \frac{1}{x} = N (f^2 n)^{2/3} \frac{1}{x}$$

$$\frac{R_E}{bN^{1/2}} = N^{1/2} (f^2 u)^{1/3} x$$

$$x^2 + \frac{1}{x} \ln \left[N^{1/2} (f^2 u)^{1/3} x \right] \stackrel{!}{=} \text{Min.}$$

$$x - \frac{1}{x^2} \ln \left[N^{1/2} (f^2 u)^{1/3} x \right] + \frac{1}{x} \frac{1}{x} = 0 \quad | \cdot x^2$$

$$x^3 = \underbrace{\ln \left[N^{1/2} (f^2 u)^{1/3} \right]} + \ln x - 1$$

N large $\Rightarrow h[\dots]$ large $\Rightarrow x$ large

$x^3 \gg h x$ $x^3 \gg \gamma$

$$\sim x \sim \left\{ h \left[N^{1/2} (f^2 u)^{1/3} \right] \right\}^{1/3}$$

$$\left. R_E \sim N b (f^2 u)^{1/3} \left\{ h \left[N^{1/2} (f^2 u)^{1/3} \right] \right\}^{1/3} \right.$$

with log-correction

"electrostatic blobs"

blob size ξ , # monomers in the blob: n

$$\xi \sim b n^{1/2}$$

electrostatic energy in one blob

$$k_B T \sim \frac{(n f e)^2}{4 \pi \epsilon \xi} \Rightarrow \Gamma \sim \frac{n^2 f^2 l_B}{\xi} \sim \frac{n^2 f^2 l_B}{b n^{1/2}}$$

$$\sim n^{3/2} f^2 l_B \Rightarrow \Gamma \sim n (f^2 l_B)^{2/3}$$

$$n \sim (f^2 u)^{-2/3}$$

$$\xi \sim b n^{1/2} \sim b (f^2 u)^{-1/3}$$

electrostatic
blob size

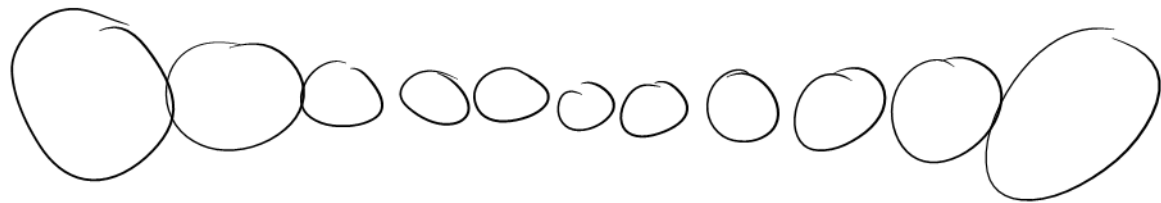
$$\# \text{ of blobs: } \frac{N}{n} \sim N (f^2 u)^{+2/3} \rightarrow$$

$$R_E \sim \frac{N}{n} \xi \sim N (f^2 u)^{2/3} b (f^2 u)^{-1/3}$$
$$\sim \left(N b (f^2 u)^{+1/3} \sim R_E \right)$$

A right answer, except log!

Simple blob picture ignores mutual
repulsion of blobs and hence
does not get the log-correction!

improvement: "inhomogeneous blobs"



much more
complicated

Dobson & Rubenstein 2005

Typically, ~~log~~ - corrections are simply
ignored!

10. Polyelectrolyte Conformations II:

Excluded-Volume Chain with Charges

blob picture: blob size ξ , n : # monomers

in a blob: $\xi \sim b n^\nu$ $\nu \approx 0.588$

$$k_B T \sim \frac{(n f e)^2}{4 \pi \epsilon \xi} \Rightarrow 1 \sim \frac{n^2 f^2 l_B}{\xi} \sim \frac{n^2 f^2 l_B}{b n^\nu}$$

$$\sim n^{2-\nu} f^2 l_B \Rightarrow 1 \sim n (f^2 l_B)^{1/(2-\nu)}$$

$$\left(n \sim (f^2 u)^{-1/(2-\nu)} \right)$$

$$\xi \sim b h^\nu \sim b (f^2 u)^{-\nu/(2-\nu)}$$

$$\# \text{ of blobs} \sim \frac{N}{n} \sim N (f^2 u)^{1/(2-\nu)}$$

$$\mathcal{R}_{\Xi} \sim \frac{N}{n} \xi \sim N (f^2 u)^{1/(2-\nu)} b (f^2 u)^{-\nu/(2-\nu)}$$

$$\left(\mathcal{R}_{\Xi} \sim N b (f^2 u)^{(1-\nu)/(2-\nu)} \right) \quad + \text{ log corrections}$$